

Duality in $2d$ (2,2)

Supersymmetric

Non-Abelian Gauge Theories

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Motivation

1991 Candelas, de la Ossa, Green, Parkes

$f(x_1, \dots, x_5)$ a Quintic polynomial

Large Volume limit

X_f : Calabi-Yau hypersurface

$$\{f(x_1, \dots, x_5) = 0\} \subset \mathbb{C}P^4$$

Singular point

Landau-Ginzburg Orbifold (Gepner model)

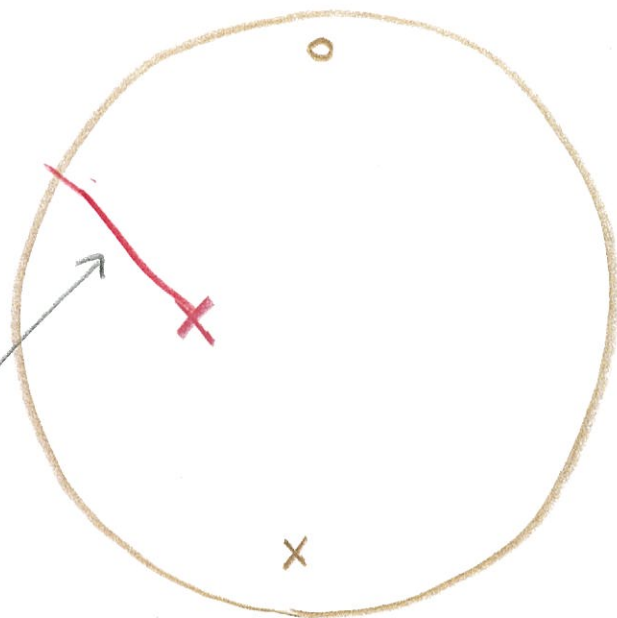
$$W = f(x_1, \dots, x_5)$$

$$\text{mod } \mathbb{Z}_5 \ni \omega : (x_1, \dots, x_5) \mapsto (\omega x_1, \dots, \omega x_5)$$

1993 Witten "Linear σ -Model"

$G = U(1), \quad P, \underbrace{X_1, \dots, X_5}_{-5 \quad 1}, \quad W = Pf(X_1, \dots, X_5)$

D-term eqn:
 $-5|P|^2 + |X|^2 = r$



\exists Coulomb branch

$r > 0$ $U(1) \xrightarrow{X \neq 0} \{1\}$

Classically \rightsquigarrow CY X_f

$r < 0$ $U(1) \xrightarrow{P \neq 0} \boxed{\mathbb{Z}_5}$

Classically \rightsquigarrow LG Orbifold

1998 Rødland

$$A_k^{ij} = -A_k^{ji} \quad i, j, k = 1, \dots, 7$$

$$X_i = (X_i^a)_{a=1,2} \quad i=1, \dots, 7$$

$$[X_i X_j] = X_i^1 X_j^2 - X_i^2 X_j^1$$

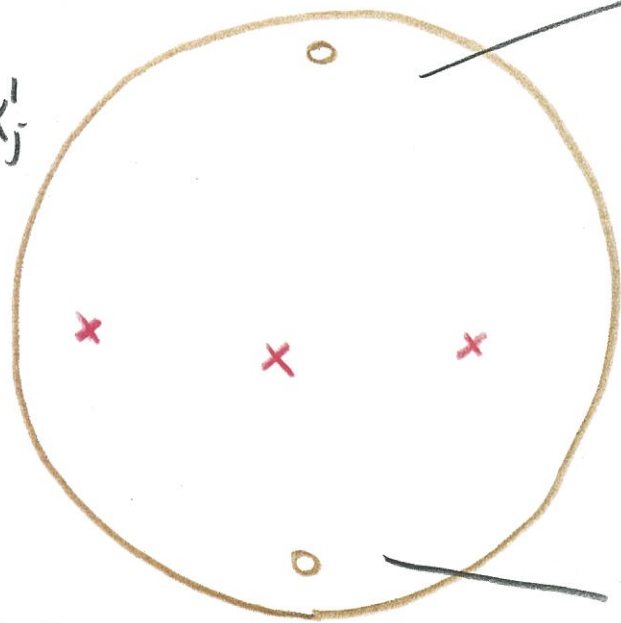
Plücker coord.

$$p^i \quad i=1, \dots, 7$$

$$A(p) = (A^{ij}(p)) \quad 7 \times 7 \text{ antisym}$$

$$A^{ij}(p) = \sum_k A_k^{ij} p^k$$

x : Singular points



X_A C.I. CY^3 in $G(2,7)$

$$= \frac{\{x \mid \sum_{ij} A_k^{ij} [X_i X_j] = 0 \quad \forall k\}}{GL(2, \mathbb{C})}$$

Y_A Pfaffian CY^3 in $\mathbb{C}P^6$

$$= \{p \mid \text{rank } A(p) \leq 4\} / \mathbb{C}^*$$

2006 KH-D.Tong

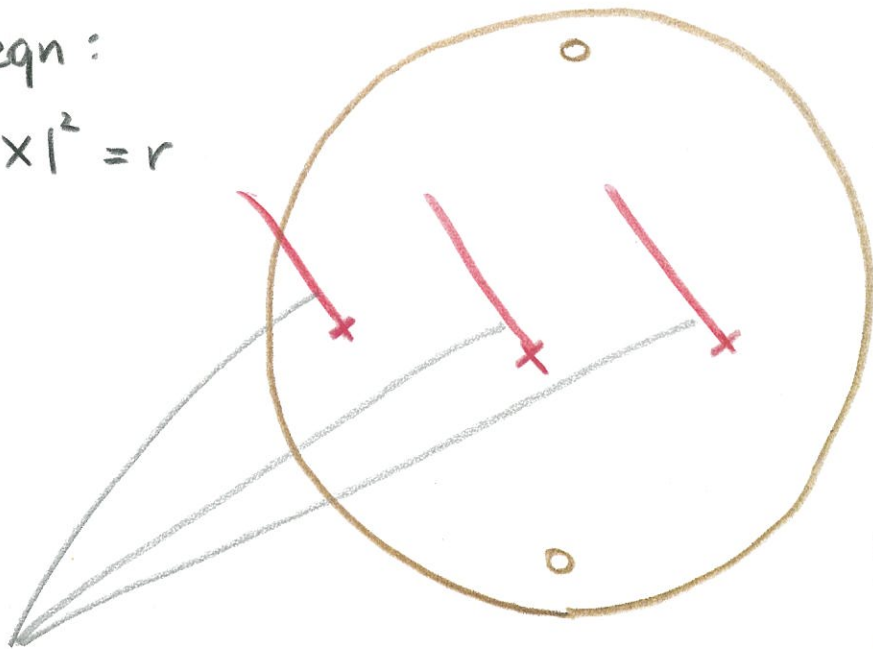
L5M

$$G = U(2), \quad \underbrace{p^1, \dots, p^7}_{\det^{-1}}, \quad \underbrace{x_1, \dots, x_7}_{\mathbb{R}}$$

$$W = \sum_{i,j,k} A_k^{ij} p^k [x_i x_j]$$

D-term eqn:

$$-|p|^2 + |x|^2 = r$$



≡ Coulomb branch

$$r > 0 : U(2) \xrightarrow{x \neq 0} \{1\}$$

Classically $\rightsquigarrow X_A$

$$r < 0 : U(2) \xrightarrow{p \neq 0} \boxed{SU(2)}$$

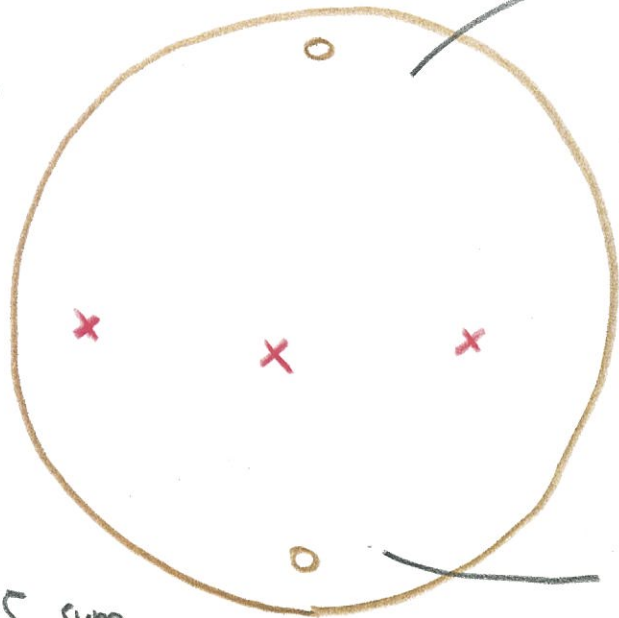
Quantum $\rightsquigarrow Y_A$

2011 Hosono-Takagi

$$S_k^{ij} = S_k^{ji} \quad i, j, k = 1, \dots, 5$$

$$X_i = (X_i^a)_{a=1,2} \quad i=1, \dots, 5$$

$$(X_i X_j) = X_i^1 X_j^1 + X_i^2 X_j^2$$



X_S Reye Congruence

(a CY³ in $\mathbb{CP}^4 \times \mathbb{CP}^4 / \mathbb{Z}_2$ exch.)

$$= \frac{\{x \mid \sum_{ij} S_k^{ij} (X_i X_j) = 0 \quad \forall k\}}{(\mathbb{C}^x \times O(2, \mathbb{C})) / \{(\pm 1, \pm 1_2)\}}$$

$$p^i = 1, \dots, 5$$

$$S(p) = (S^{ij}(p)) \quad 5 \times 5 \text{ sym.}$$

$$S^{ij}(p) = \sum_k S_k^{ij} p^k$$

x: Singular points

Y_S

ramified double cover

$\downarrow \mathbb{Z}_2$

$$Y_S = \{p \mid \text{rank } S(p) \leq 4\} / \mathbb{C}^x \subset \mathbb{CP}^4$$

\cup

ram. along $C_S = \{p \mid \text{rank } S(p) \leq 3\} / \mathbb{C}^x$

2011 KH

L & M

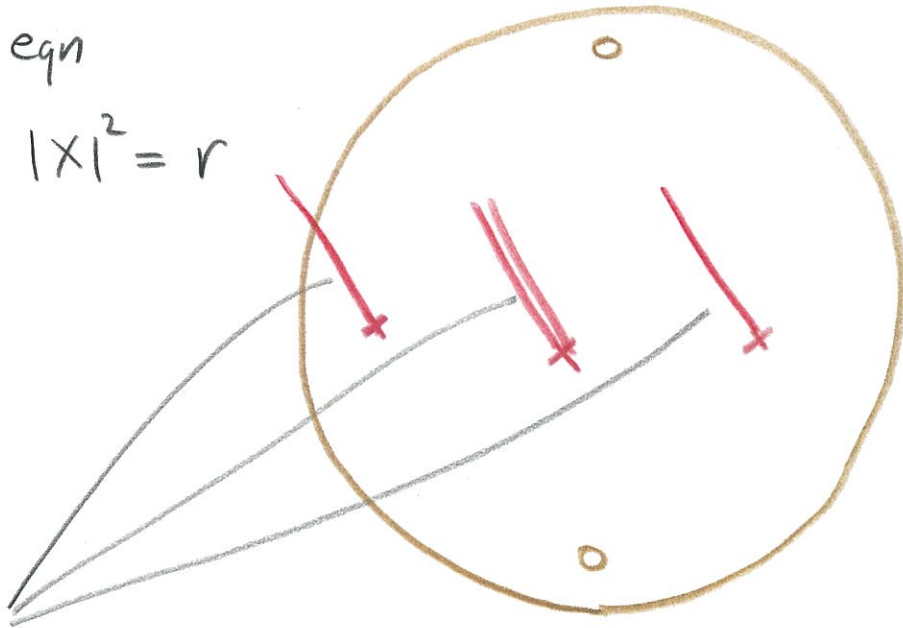
$$G = \frac{U(1) \times O(2)}{\{(\pm 1, \pm 1)\}}$$

$$\underbrace{P^1, \dots, P^5}_{(-2, 1)}, \quad \underbrace{X_1, \dots, X_5}_{(1, 2)}$$

$$W = \sum_{i,j,k} S_{ij}^k P^k(X_i, X_j)$$

D-term eqn

$$-2|P|^2 + |X|^2 = r$$



\exists Coulomb branch

$$r > 0 : G \xrightarrow{X \neq 0} \{1\}$$

Classically
→

X_5

$$r < 0 : G \xrightarrow{P \neq 0} O(2)$$

Quantum
→

$\sim Y_5$

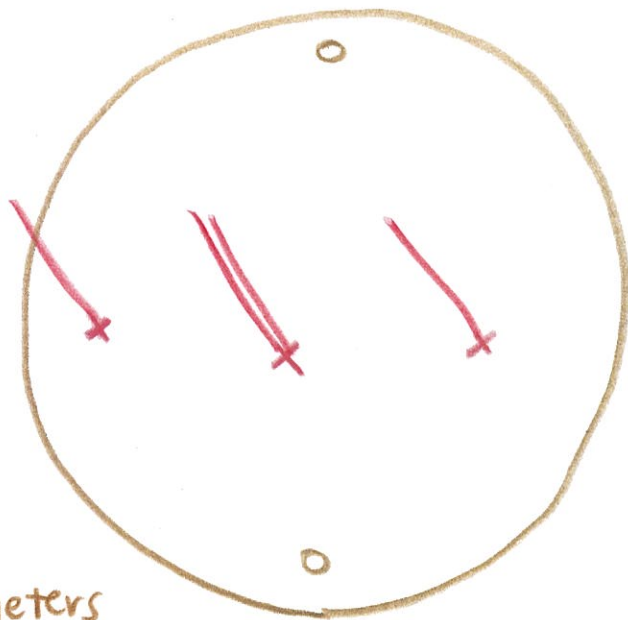
\exists Dual Linear σ -Model (for Hosono-Takagi)

$$\tilde{G} = \frac{U(1) \times SO(4)}{\{(\pm 1, \pm 1_4)\}}, \quad \underbrace{P^1, \dots, P^5}_{(-2, 1)}, \quad \underbrace{\tilde{X}^1, \dots, \tilde{X}^5}_{(-1, 4)}, \quad \underbrace{S_{ij}}_{(1, 1)}^{S_{ji}}$$

$$W = \sum S_{ij} (\tilde{X}^i \tilde{X}^j) + \sum S_k^{ij} P^k S_{ij}$$

D-term eqn

$$-2|P|^2 - |\tilde{X}|^2 + |S|^2 = \tilde{r}$$



$$\tilde{r} > 0: \tilde{G} \xrightarrow{S \neq 0} SO(4)$$

Quantum \rightarrow

X_S

$$\tilde{r} < 0: \tilde{G} \xrightarrow{\substack{\tilde{X} \neq 0 \\ P \neq 0}} \boxed{\{1\}}$$

Classically \rightarrow

\tilde{Y}_S

Relation of parameters

$$\begin{cases} r = \tilde{r} + 5 \log 2 \\ \theta \equiv \tilde{\theta} + \pi \end{cases}$$

Dual Linear σ -Model (for Rødland)

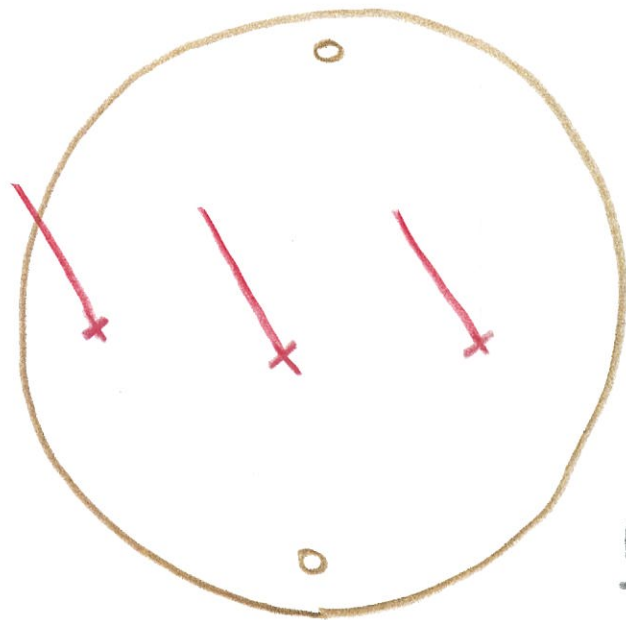
$$\tilde{G} = \frac{U(1) \times USp(4)}{\{(\pm 1, \pm 1_4)\}}$$

$\underbrace{P^1, \dots, P^7}_{(-2, 1)}, \underbrace{\tilde{X}^1, \dots, \tilde{X}^7}_{(-1, 4)}, \underbrace{a_{ij}}_{(1, 1)} \stackrel{-a_{ji}}{=}$

$$W = \sum a_{ij} [\tilde{X}^i \tilde{X}^j] + \sum A_k^{ij} P^k a_{ij}$$

D-term eqn

$$-2|p|^2 - |\tilde{X}|^2 + |a|^2 = \tilde{r}$$



Relation

$$\begin{cases} r = \tilde{r} \\ \theta = \tilde{\theta} \end{cases}$$

$$\tilde{r} > 0 : \tilde{G} \xrightarrow{a \neq 0} USp(4)$$

Quantum X_A

$$\tilde{r} < 0 : \tilde{G} \xrightarrow{\substack{\tilde{X} \neq 0 \\ P \neq 0}} \boxed{\{1\}}$$

Classically Y_A

$M_K \ni x, y \Rightarrow D_B(x) \cong D_B(y)$

Equivalence of
Categories of B-branes

$D^b(X_f) \cong MF_{\mathbb{Z}_5}(f)$

Orlov 2005

$D^b(X_A) \cong D^b(Y_A)$

Borisov-Caldararu, Kuznetsov 2006

$D^b(X_S) \cong D^b(\tilde{Y}_S)$

Hosono-Takagi 2011

Note $X_A \not\cong Y_A$, $X_S \not\cong \tilde{Y}_S$ Birationally inequivalent

The Models

$$k = 1, 2, 3, \dots, \quad N = 0, 1, 2, \dots$$

Gauge group $G = O(k)$ or $SO(k)$

Matters : N fields in \mathbb{C}^k X_1, \dots, X_N

Superpotential : $W = 0$

• $O(k) = SO(k) \rtimes \mathbb{Z}_2 \dots \mathbb{Z}_2$ orbifold of $SO(k)$ theory

2 possibilities $O_+(k), O_-(k)$.

• $k \geq 3$: $\pi_1(G) \cong \mathbb{Z}_2 \rightarrow$ mod 2 θ -angle , No F.I.

• $k=2$: $SO(2)$: $\theta \in \mathbb{R}/2\pi\mathbb{Z}$, F.I. $\in \mathbb{R}$; $O_{\pm}(2)$: $\theta = 0$ or π , No F.I.

$$k = 2, 4, 6, \dots, \quad N = 0, 1, 2, \dots$$

Gauge group $G = USp(k)$

Matters: N fields in \mathbb{C}^k x_1, \dots, x_N

Superpotential : $W = 0$

$$\pi_1(G) = \{0\} \quad \rightarrow \quad \text{No } \theta, \text{ No F.I.}$$

Regularity

To find the nature of low energy theory,

Let's first look at $\{U=0\}$.

• scalar fields : σ in vector, X in chiral

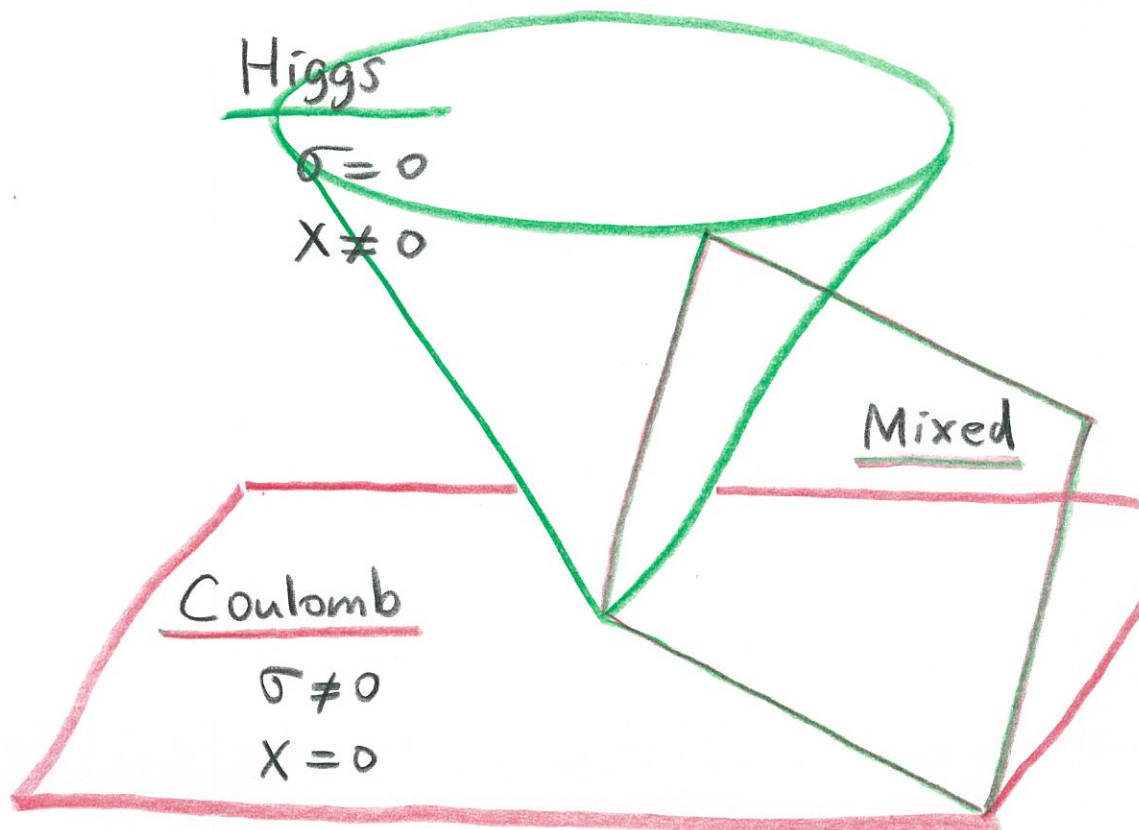
• Potential $U(\sigma, X) = \frac{1}{2g^2} |[\sigma, \sigma^\dagger]|^2 + \frac{1}{2} |\sigma X|^2 + \frac{1}{2} |\sigma^\dagger X|^2 + \frac{g^2}{2} |\mu(X)|^2$

$U=0$: each term = 0

$[\sigma, \sigma^\dagger] = 0, \sigma X = \sigma^\dagger X = 0, \mu(X) = 0.$

↑
e.g. $O(k)$ $k=2l$

$$\sigma = \begin{pmatrix} 0 & -\sigma_1 & & & & \\ \sigma_1 & 0 & & & & \\ & & \ddots & & & \\ & & & 0 & -\sigma_l & \\ & & & \sigma_l & 0 & \end{pmatrix}$$



Problem: Non-compactness.

Higgs branch may be lifted by turning on $W(X)$,
but that does not change Coulomb branch.

In fact, \exists quantum correction to U on the Coulomb.

$$\sigma \neq 0 : G \longrightarrow U(1)^k$$

Copies of Maxwell theory

$$\Delta U = \frac{g^2}{2} |\theta_{\text{eff}}|^2$$

Electro-Static Energy

$$G = O(k) \text{ or } SO(k) : \theta_{\text{eff}} \equiv \pi(N-k) + \theta_0 \pmod{2\pi}$$

$$\left(\begin{array}{l} \theta_0 = 0 \text{ or } \pi : \\ \text{mod } 2 \text{ } \theta\text{-angle} \end{array} \right)$$

$$G = USp(k) : \theta_{\text{eff}} \equiv \pi(N-k) \pmod{2\pi}$$

Thus, the Coulomb and the mixed branches are lifted by quantum correction if

$$O(k), SO(k) : \quad N-k \text{ odd, } \theta_0 \equiv 0$$

$$N-k \text{ even, } \theta_0 \equiv \pi$$

$$USp(k) : \quad N-k \text{ odd.}$$

Such theories are called regular.

Results

$G = O_{\pm}(k), SO(k)$, X_1, \dots, X_N in \mathbb{C}^k , $W=0$

• $N \leq k-2$: SUSY

↓ regular

• $N = k-1$: Free theory of mesons $(X_i X_j) = \sum_{a=1}^k X_i^a X_j^a$
 $\mathbb{C}^{\frac{N(N+1)}{2}}$ for $SO(k), O_-(k)$, $\mathbb{C}^{\frac{N(N+1)}{2}} \sqcup \mathbb{C}^{\frac{N(N+1)}{2}}$ for $O_+(k)$.

• $N \geq k$: \exists duality

$$G = O_+(k) \iff \tilde{G} = SO(N-k+1)$$

$$SO(k) \iff O_+(N-k+1)$$

$$O_-(k) \iff O_-(N-k+1)$$

The dual theory has N fields $\tilde{X}^1, \dots, \tilde{X}^N$ in \mathbb{C}^{N-k+1}

$$\frac{N(N+1)}{2} \text{ singlets } S_{ij} = S_{j\bar{i}}$$

$$W = \sum_{i,j} S_{ij} (\tilde{X}^i \tilde{X}^j)$$

• mesons $(x_i x_j)$ \longleftrightarrow singlets S_{ij}

• $G = SO(k)$ vs $\tilde{G} = O_+(N-k+1)$

\mathbb{Z}_2 global symmetry \longleftrightarrow Quantum \mathbb{Z}_2 symmetry
 $O(k)/SO(k)$ of the \mathbb{Z}_2 orbifold

baryons $[x_{i_1} \dots x_{i_k}]$ \longleftrightarrow twist operators

• $G = O_+(k)$ vs $\tilde{G} = SO(N-k+1)$

Similar

• $G = O_-(k)$ vs $\tilde{G} = O_-(N-k+1)$

Quantum \mathbb{Z}_2 symmetry \longleftrightarrow quantum $\mathbb{Z}_2 (-1)^F$ symmetry.

$G = USp(k), X_1, \dots, X_N \text{ in } \mathbb{C}^k, W=0$

• $N \leq k$: ~~SUSY~~

↓ regular • $N = k+1$: Free theory of mesons $[X_i X_j] = \sum_{a,b} X_i^a J_{ab} X_j^b$
ie. CFT of $\mathbb{C}^{\frac{N(N-1)}{2}}$.

• $N \geq k+3$: \exists dual theory

$$\tilde{G} = USp(N-k-1)$$

N fields $\tilde{X}^1, \dots, \tilde{X}^N$ in \mathbb{C}^{N-k-1}

$\frac{N(N-1)}{2}$ singlets $a_{ij} = -a_{ji}$

$$W = \sum_{i,j} a_{ij} [\tilde{X}^i \tilde{X}^j]$$

mesons $[X_i X_j] \longleftrightarrow$ singlets a_{ij}

Tests of Duality

o Check at low values of k, N (cf. $O(1) = \mathbb{Z}_2, SO(1) = \{1\}$)

o Symmetry

O, SO

	$SU(N)$	$U(1)_B$	$U(1)_V$	$U(1)_A$
X	\square	1	0	0
\tilde{X}	$\bar{\square}$	-1	1	0
S	\boxplus	2	0	0

$$\hat{C} = kN - \frac{k(k-1)}{2}$$

match!

$$\hat{C} = \frac{N(N+1)}{2} - \frac{(N-k+1)(N-k)}{2}$$

USp

	$SU(N)$	$U(1)_B$	$U(1)_V$	$U(1)_A$
X	\square	1	0	0
\tilde{X}	$\bar{\square}$	-1	1	0
a	\boxplus	2	0	0

$$\hat{C} = kN - \frac{k(k+1)}{2}$$

match!

$$\hat{C} = \frac{N(N-1)}{2} - \frac{(N-k-1)(N-k)}{2}$$

o Born-Oppenheimer for dual.

$$\left. \begin{array}{l} \text{corank}(S) \leq (N-k+1)-2 : \text{SUSY} \\ \text{corank}(S) = (N-k+1)-1 : \text{Free theory of mesons} \end{array} \right\} \Rightarrow \underline{\text{rank}(S) \leq k}$$

$$\left. \begin{array}{l} \text{corank}(a) \leq (N-k-1) : \text{SUSY} \\ \text{corank}(a) = (N-k-1)+1 : \text{Free theory of mesons} \end{array} \right\} \Rightarrow \underline{\text{rank}(a) \leq k}$$

o Flow in Complex mass

$$\begin{array}{ccc} O_+(k), N & \longleftrightarrow & SO(N-k+1), N \\ \Delta W = m(X_N X_N) \downarrow & & \downarrow \Delta W = m S_{NN} \rightsquigarrow (\tilde{X}^N \tilde{X}^N) + m = 0 \\ O_+(k), N-1 & & SO(N-k), N-1 \end{array}$$

o Vacuum Counting with twisted masses. ✓

o Comparison of chiral rings : (c,c) & (a,c) . ✓