

①

Towards A Global Mirror Symmetry

YONGBIN RUAN

University of Michigan

Plan of talk:

- (I) What is "global" ?
- (II) A conjectural physical package
 \Rightarrow
BCOV
- (III) An approach to hypersurface
LG -model
- (IV) Un-expected bonus ! \Downarrow
orbifold Gw -Theory

(I) WHAT IS "GLOBAL"?

(2)

"Local" Mirror
symmetry



Mirror symmetry of
local Calabi-Yau

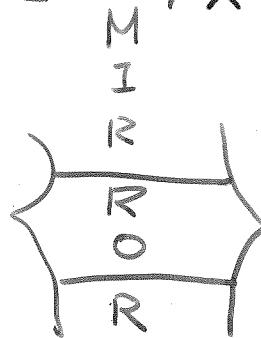
X - CY 3-fold $\longleftrightarrow X^\vee$ - another CY 3-fold

A-model



Kahler str

GW- Theory
($g=0$)



B-model



complex str

periods

Famous
example:

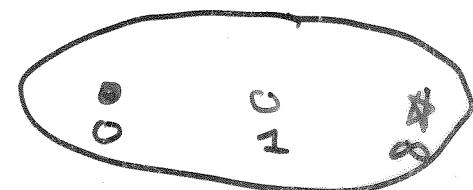
$$X = \left\{ \sum_{i=1}^5 x_i^5 = 0 \right\} \longleftrightarrow X^\vee = \left\{ \sum_{i=1}^5 x_i^5 - 5 \sum_{i=1}^5 x_i = 0 \right\}$$

~~\mathbb{Z}_5^3~~

Kahler str

$t \in$ Kahler Cone $\xleftarrow[\text{match } \chi \in]{\text{NO}} \text{complex str}$

$$\pi_1 = 1$$



$$\mathbb{P}^1 - \{0, 1, \infty\}$$

$$\pi_1 \neq 1$$

A temporary
solution at
the time (20 years
ago) ↓

"local mirror" //
symmetry

Kahler
str

"Local"
 $t \longleftrightarrow \gamma \in \{ \text{a neigh of} \}$

$\gamma = \infty$
↑
large complex
str limit

Now: Restore "Global" structure of Mirror
Symmetry

Benefit: { (1) compute higher genus Gromov-Witten
Theory
(2) study modularity of ..
(3) Prove Landau-Ginzburg /
Calabi-Yau correspondence.
(4) More, ...

"GLOBAL" = allow γ to move
around in the entire
moduli space of
complex str

(II) A conjectural physical package

(4)

Background: B-model ($g=0$)
ss
Periods

A toy model

elliptic curve

$$E(\chi) = \left\{ \sum_{i=1}^3 x_i^3 - 3\chi^{\frac{1}{3}} \prod_{i=1}^3 x_i = 0 \right\} / z_3^2$$

$\gamma \in \mathbb{P}^1 - \{0, 1, \infty\}$

CY-form:

$$H^{1,0}(E(\chi)) = \mathbb{C}$$

ω_χ - holomorphic (1,0) form

vary holomorphically as we vary χ

another
choice.

$$\omega(\chi) \longrightarrow f(\chi) \omega(\chi)$$

holomorph.
funct

(5)

$H_1(E|\mathcal{W}, \mathbb{Z})$ has a symplectic basis

$$A, B \text{ with } A^2 = B^2 = 0, A \cdot B = 1, B \cdot A = -1$$

change
basis

$$\begin{pmatrix} A \\ B \end{pmatrix} \rightarrow \begin{pmatrix} A' \\ B' \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

π

$$SL_2 \mathbb{Z}$$

PERIOD:

$$q = \int_A \omega(x) \quad p = \int_B \omega(x)$$

Vary $\chi \rightarrow q(\chi), p(\chi)$ - multi-valued
funct of $\mathbb{P}' - \{0, 1, \infty\}$.

$$\tau = \frac{p(\chi)}{q(\chi)} \in \mathbb{H}_+ - \text{upper half-plane}$$

\downarrow
 χ
 $\mathbb{P}' - \{0, 1, \infty\}$ - universal cover

$$\mathbb{H}_+ / \begin{matrix} \text{monodromy} \\ \text{group} \end{matrix} = P(3) \subset SL_2 \mathbb{Z}$$

GW-theory \iff
of kahler
parameter
 t

Periods of
 χ near $\{\infty\}$ with a special choice
of basis A, B

$\tau = i\infty$ "large complex str.
limit"

(6)

Going Back to quintic 3-fold (B-model)

$H_+ \rightarrow \mathbb{C}$ -modular coord

$$\gamma \in \text{RP}^1 - \{0, 1, \infty\} = \frac{H_+}{P} \subset \text{Sp}(2, \mathbb{Z})$$

$$H_3(X^\vee, \mathbb{Z}) = 4$$

GW-theory
of quintic
 $g=0$

"local"
↔
Mirror

Periods near $\tau = i\infty$
large complex limit

Remark

(1) Above "local" mirror symmetry for $g=0$
has been verified by Givental, Liu-Liang-Yau
in the middle of 90's.

Our interest

- (2) Higher genus ($g > 0$) GW-theory
- (3) "Global" properties of GW-function.

A conjectural physical package (BCOV) (7)

(I) GLOBAL PROPERTY:

- ① exist a global $\mathcal{F}_g^B(\tau, \bar{\tau})$ - genus g generating funct of B-model GW Theory
- (i) defined for all τ
 - (ii) non-holomorphic

(2) Modular Invariance

$$\mathcal{F}_g^B(h\tau, h\bar{\tau}) = \mathcal{F}_g^B(\tau, \bar{\tau}) j(h, \tau)^k$$

$h \in P$ - monodromy group

- ③ $\partial_{\bar{\tau}} \mathcal{F}_g^B = P_g (\mathcal{F}_{g+g}^B)$ - holomorphic anomaly equation
- ↑ BCOV, Klemm,

$$\mathcal{F}_g^B(\tau, \bar{\tau}) = \sum_j \mathcal{F}_{g,j}^B(\tau) (im\tau)^{-j}$$

- (2) + (3) $\Rightarrow \mathcal{F}_g^B(\tau)$ - quasi-modular form

(II) SPECIAL LIMIT (Mirror Conjectures)

B-model

CY-to-CY

A-model

(1) near $T = i\infty$ Gromov-Witten
theory of X

Original "local" mirror

Small / Large
duality

Symmetry

= Landau Ginzburg / CY
correspondence(2) near $\gamma = 0$ \longleftrightarrow "conjectural"
(Gepner limit) CY-to-LG Landau-Ginzburg
model(3) near $\gamma = 1$ \longleftrightarrow "conjectural"
(conifold limit) matrix model "

(4) Beyond quintic, other limits.

Klemm's group:Assume the existence of
above packageknown for mathematician
only for geo. 1

+ general properties of

special limits

A STRIKING computation of Gw-theory of,
 $g \leq s_1$

(8')

Goal of Remaining talk

- Describe an approach for a mathematical construction of above package for hypersurface.
Most of steps are conjectural at this moment
- Present some theorems in dimension one.

(9)

First advance: Gepner limit

(Conjectural LG-model = Theory of
Fan-Jarvis-Ruan-Witten
(2007))

LG - model:

- (i) $W: \mathbb{C}^N \rightarrow \mathbb{C}$ "non-degenerate"
quasi-homogeneous
poly
- (ii) $G \subset \text{Aut}(W)$ - finite abelian
symmetry

Theory of Fan-Jarvis-Ruan-Witten:

- A complete A-model theory of LG-model based on solving Witten eqn

$$\bar{\partial} s_i + \overline{\partial_i W} = 0$$
- A GW-type curve counting theory
 - based on $\overline{\mathcal{M}}_{g,n}$
 - 2D TFT
 - satisfies axioms of GW-theory
- Much easier to calculate (ADE, elliptic singularity $g=0$ quintic, expanding quickly)

(10)

what happen for (I): build a rigorous theory of $F_g^B(\tau, \bar{\epsilon})$ with expected properties

- A hard problem
- Progress on $X = T^{2n}$ (Costello - Li)
- Our approach for hypersurfaces such as quintic 3-fold (Milanov - Krautz - Shen)

Key observation:

γ -deformation: $x_1^3 + x_2^3 + x_3^3 - 3\gamma^{\frac{1}{3}}x_1 x_2 x_3$

subset

$$x_1^3 + x_2^3 + x_3^3 - 3\gamma^{\frac{1}{3}}x_1 x_2 x_3 + t_0 + t, x,$$

$$+ t_2 x_2 + t_3 x_3 + t_4 x_1 x_2 + t_5 x_1 x_3 + t_6 x_2 x_3$$

"miniversal" deformation of singularity

$$W = x_1^3 + x_2^3 + x_3^3$$

A "baby" B-model of singularity /LG-model

$$W = x_1^3 + x_2^3 + x_3^3$$

Milnor ring $Q_W = \frac{\mathbb{C}[x_1, x_2, x_3]}{\partial_{x_1} W, \partial_{x_2} W, \partial_{x_3} W}$ 8-dim
 $= \{1, x_1, x_2, x_3, x_1 x_2, x_1 x_3, x_2 x_3, x_1 x_2 x_3\}$

Miniversal deformation $W_{(t_i, \sigma)} = x_1^3 + x_2^3 + x_3^3 + \sigma x_1 x_2 x_3 + t_0 + t_1 x_1 + t_2 x_2 + t_3 x_3 + t_4 x_1 x_2 + t_5 x_1 x_3 + t_6 x_2 x_3$

B-model parameter space $= \{(t_i, \sigma), |t_i| < \epsilon, \sigma^3 \neq 27\}$

For each $\{t_i, \sigma\}$.

$$Q_{W_{(t_i, \sigma)}} = \frac{\mathbb{C}[x_1, x_2, x_3]}{\partial_{x_1} W_{(t_i, \sigma)}, \partial_{x_2} W_{(t_i, \sigma)}, \partial_{x_3} W_{(t_i, \sigma)}}$$

a family of Frobenius algebra

Pairing: $\langle \phi_1, \phi_2 \rangle = \text{Res} \frac{\phi_1 \phi_2 dx_1 dx_2 dx_3}{dW_{(t_i, \sigma)}}$

Main properties

① For generic $t_i \neq 0$, $W_{(t_i, \sigma)}$ - holomorphic

Morse funct
↓

Frobenius algebra
of such (t_i, σ) is semi-simple

② pairing is NOT flat

Saito - Givental Theory

(I) Saito - Theory:

replace $dx_1 dx_2 dx_3 \rightarrow \frac{1}{P(\sigma)} dx_1 dx_2 dx_3$ primitive form
↓

flat pairing \Rightarrow Frobenius manifold str

$F_o^B(t_i, \sigma, P)$
 $(t_i \neq 0)$ - semi-simple

(II) Givental Theory:

on semi-simple Frobenius model, exist

$F_g^B(t_i, \sigma, P)$ $(t_i \neq 0)$

what we know about primitive form

(13)

$$\frac{1}{P} d\bar{x} \quad ?$$

-
- P always exist locally \Rightarrow local Saito-Givental Theory
 - explicit formula for ADE, elliptic singularities
 - difficult to get explicit formula in general
 - along cy-direction (marginal deformation), related to periods.

↓ leads to

Global Saito - Givental Theory

(under developed by Milanov, —)

(14)

Global Saito-Givental Theory in $\dim = 1$ (Milanov, -)

starting point:

$$P(\sigma) - \text{period, i.e } P(\sigma) = \int_A w(\sigma)$$

$$\text{Recall } \tau = \frac{\int_B w(\sigma)}{\int_A w(\sigma)} \in \mathbb{H}_+$$

A, B - symplectic basis

flat coord along σ -direction

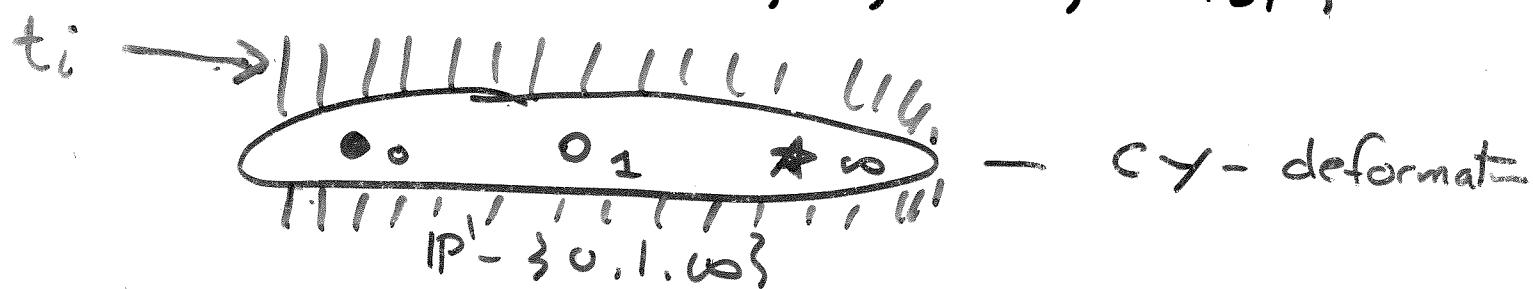
B-model

$$\text{Parameter} = \{(t_i, \tau), |t_i| < \varepsilon, \tau \in \mathbb{H}_+\}.$$

Space

↓ / monodromy group

$$\{(t_i, \sigma), |t_i| < \varepsilon, \sigma^3 \neq 27\}$$



$$(t_i, \tau) \rightarrow \text{Frobenius mod str} \rightarrow \mathcal{F}_g^B(t_i, \tau)$$

$$(t_i \neq 0, \tau) \rightarrow \mathcal{F}_g^B(t_i, \tau) \xrightarrow{\text{holomorphic}}$$

semi-simple

with respect to t_i, τ

Non-Modular
invariant
(Milanov.)

Let

$$\mathcal{D}_{SG}^B(t_i, \tau) = \exp\left(\sum_{g \geq 0} t_i^{2g-2} \mathcal{F}_g^B\right)$$

$h \in P$

$$\mathcal{D}_{SG}^B(t_i; h\tau) = \hat{X}_h \mathcal{D}_{SG}^B$$

where \hat{X}_h - differential operator
defined out of $h \in \Gamma(3)$

Corollary: $\mathcal{F}_g^B(t_i, \tau)$ is not modular

A magic trick: Anti-holomorphic completion

We found an explicit way to complete

$$\mathcal{F}_g^B(t_i, \tau) \rightarrow \mathcal{F}_g^B(t_i, \tau, \bar{\tau})$$

$$\mathcal{F}_g^B(t_i, \tau, \bar{\tau}) = \sum_{j=1}^k \mathcal{F}_{g,j}^B(\tau) (\text{im } \tau)^{-j}$$



defined via Feynman diagram
expansion

quasi-modular
form

motivated by Aganagic-Bouchard-Klemm

Easy Fact: $\mathcal{F}_g^B(t_i, \tau, \bar{\tau})$ satisfies holomorphic anomaly equation

Modular

Invariance : $\mathcal{F}_g^B(t_i, \tau, \bar{\tau})$ is modular
(Milanov, -)

Assume: $\mathcal{F}_g^B(t_i \neq 0, \tau, \bar{\tau})$ extends to $t_i = 0$

↑

$$\mathcal{F}_g^B(\tau, \bar{\tau}) = \mathcal{F}_g^B(0, \tau, \bar{\tau})$$

somehow a difficult problem!

Corollary (a):

$$\mathcal{F}_g^B(t_i, \tau, \bar{\tau}) = \sum_I a_I(\tau, \bar{\tau}) t_i$$

then $a_I(\tau, \bar{\tau})$ are classical modular forms almost holomorphic

$$\text{or } \mathcal{F}_g^B(t_i, \tau) = \sum_I a_I(\tau) t_i$$

classical quasi-modular form

Corollary (b)

$\mathcal{F}_g^B(\tau, \bar{\tau})$ satisfies holomorphic anomaly equation
desired B -model theory

(IV) An un-expected Bonus !

(17)

!!

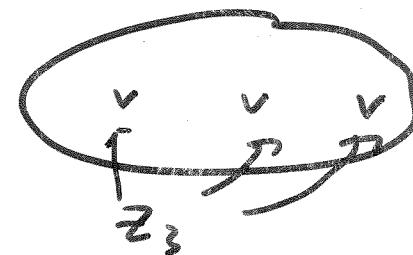
$\mathcal{F}_g^B(t_i, \tau)$ has a mirror of !
its own .

A-model $\quad \{x_1^3 + x_2^3 + x_3^3 = 0\} / z_3^2$

Mirror :



$$\mathbb{P}_{3,3,3}^1 \stackrel{!!}{=} \text{orbifold } \mathbb{P}^1$$



LG-dual ; ($W = x_1^3 + x_2^3 + x_3^3, z_3^3$)

Theorem : (Krawitz - Shen)

① Near $\tau = 0$, $\mathcal{F}_g^B(t_i, \tau) = \mathcal{F}_g^{FJRW}(t'_i, \tau')$

\downarrow
extends to \Downarrow extends to $t'_i = 0$

② Near $\tau = i\infty$, $\mathcal{F}_g^B(t_i, \tau) = \mathcal{F}_g^{GW}(t'_i, q = e^{\frac{2\pi i \tau}{3}})$

③ Same holds for $x_9 = x_1^2 + x_2^4 + x_3^4 \Leftrightarrow \mathbb{P}_{2,4,4}^1$
 $J_{10} = x_1^2 + x_2^3 + x_3^6 \Leftrightarrow \mathbb{P}_{2,3,6}^1$

Two Bonuses

(I) GW-Theories of $P_{3,3,3}^1, P_{2,4,4}^1$

$P_{2,3,6}^1$ are quasi-modular



wanted very much by mathematician

(II) LG/CY - correspondence

holds for all general for

these examples.



First example of all genera

A less important

Result :

Restrict to

$$t_i = 0$$

\implies recover elliptic curve

Known already by Okounkov
Pandharipande