Math 501 - Differential Geometry Herman Gluck
Tuesday March 13, 2012

## 6. GEODESICS

In the Euclidean plane, a straight line can be characterized in two different ways:
(1) it is the shortest path between any two points on it;
(2) it bends neither to the left nor the right (that is, it has zero curvature) as you travel along it.

We will transfer these ideas to a regular surface in 3-space, where geodesics play the role of straight lines.

## Covariant derivatives.

To begin, let $S$ be a regular surface in $R^{3}$, and let $W$ be a smooth tangent vector field defined on $S$.


If $p$ is a point of $S$ and $Y$ is a tangent vector to $S$ at $p$, that is, $Y \in T_{p} S$, we want to figure out how to measure the rate of change of $W$ at $p$ with respect to $Y$.

Let $\alpha(t)$ be a smooth curve on $S$ defined for $t$ in some neighborhood of 0 , with $\alpha(0)=\mathrm{p}$, and $\alpha^{\prime}(0)=\mathrm{Y}$.

Then $\mathrm{W}(\alpha(\mathrm{t}))=\mathrm{W}(\mathrm{t})$ is a vector field along the curve $\alpha$.


We define

$$
(\mathrm{DW} / \mathrm{dt})(\mathrm{p})=\text { orthog proj of } \mathrm{dW} /\left.\mathrm{dt}\right|_{\mathrm{t}=0} \text { onto } \mathrm{T}_{\mathrm{p}} \mathrm{~S}
$$

and call this the covariant derivative of the vector field W at the point p with respect to the vector Y .

The above definition makes use of the extrinsic geometry of $S$ by taking the ordinary derivative $d W / d t$ in $R^{3}$, and then projecting it onto the tangent plane to $S$ at $p$.

But we will see that, in spite of appearances, the covariant derivative $\mathrm{DW} / \mathrm{dt}$ depends only on the intrinsic geometry of $S$.

To show that the covariant derivative depends only on the intrinsic geometry of $S$, and also that it depends only on the tangent vector $Y$ ( not the curve $\alpha$ ), we will obtain a formula for $D W / d t$ in terms of a parametrization $X(u, v)$ of $S$ near $p$.

Let $\alpha(\mathrm{t})=\mathrm{X}(\mathrm{u}(\mathrm{t}), \mathrm{v}(\mathrm{t}))$, and write

$$
\begin{aligned}
\mathrm{W}(\mathrm{t}) & =\mathrm{a}(\mathrm{u}(\mathrm{t}), \mathrm{v}(\mathrm{t})) \mathrm{X}_{\mathrm{u}}+\mathrm{b}(\mathrm{u}(\mathrm{t}), \mathrm{v}(\mathrm{t})) \mathrm{X}_{\mathrm{v}} \\
& =\mathrm{a}(\mathrm{t}) \mathrm{X}_{\mathrm{u}}+\mathrm{b}(\mathrm{t}) \mathrm{X}_{\mathrm{v}} .
\end{aligned}
$$

Then by the chain rule,

$$
\begin{aligned}
\mathrm{dW} / \mathrm{dt}=\mathrm{W}^{\prime}(\mathrm{t})= & \mathrm{a}^{\prime} \mathrm{X}_{\mathrm{u}}+\mathrm{a}\left(\mathrm{X}_{\mathrm{u}}\right)^{\prime}+\mathrm{b}^{\prime} \mathrm{X}_{\mathrm{v}}+\mathrm{b}\left(\mathrm{X}_{\mathrm{v}}\right)^{\prime} \\
= & \mathrm{a}^{\prime} \mathrm{X}_{\mathrm{u}}+\mathrm{a}\left(\mathrm{X}_{\mathrm{uu}} \mathrm{u}^{\prime}+\mathrm{X}_{\mathrm{uv}} \mathrm{v}^{\prime}\right) \\
& +\mathrm{b}^{\prime} \mathrm{X}_{\mathrm{v}}+\mathrm{b}\left(\mathrm{X}_{\mathrm{vu}} \mathrm{u}^{\prime}+\mathrm{X}_{\mathrm{vv}} \mathrm{v}^{\prime}\right) .
\end{aligned}
$$

## Recall that

$$
\begin{aligned}
& \mathrm{X}_{\mathrm{uu}}=\Gamma^{1}{ }_{11} \mathrm{X}_{\mathrm{u}}+\Gamma_{11}^{2} \mathrm{X}_{\mathrm{v}}+\mathrm{e} \mathrm{~N} \\
& \mathrm{X}_{\mathrm{uv}}=\Gamma_{12}^{1} \mathrm{X}_{\mathrm{u}}+\Gamma_{12}^{2} \mathrm{X}_{\mathrm{v}}+\mathrm{f} \mathrm{~N} \\
& \mathrm{X}_{\mathrm{vu}}=\Gamma^{1}{ }_{21} \mathrm{X}_{\mathrm{u}}+\Gamma_{21}^{2} \mathrm{X}_{\mathrm{v}}+\mathrm{f} \mathrm{~N} \\
& \mathrm{X}_{\mathrm{vv}}=\Gamma_{22}^{1} \mathrm{X}_{\mathrm{u}}+\Gamma_{22}^{2} \mathrm{X}_{\mathrm{v}}+\mathrm{gN} .
\end{aligned}
$$

Inserting these values into the formula for $d W / d t$ and dropping each appearance of N , we get

DW/dt $=\quad\left(\mathrm{a}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{u}}$

$$
+\left(\mathrm{b}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{v}} .
$$

We repeat the formula:

$$
\begin{aligned}
\mathrm{DW} / \mathrm{dt}= & \left(\mathrm{a}^{\prime}+\mathrm{a} \Gamma_{11}^{1} \mathrm{u}^{\prime}+\mathrm{a} \Gamma_{12}^{1} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{u}} \\
& +\left(\mathrm{b}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{22} \mathrm{v}^{\prime}\right) X_{v} .
\end{aligned}
$$

From this formula, we learn two things:
(1) The covariant derivative $\mathrm{DW} / \mathrm{dt}$ depends only on the tangent vector $Y=X_{u} u^{\prime}+X_{v} v^{\prime}$ and not on the specific curve $\alpha$ used to "represent" it.
(2) The covariant derivative DW/dt depends only on the intrinsic geometry of the surface $S$, because the Christoffel symbols $\Gamma^{\mathrm{k}}{ }_{\mathrm{ij}}$ are already known to be intrinsic.

## Tensor notation.

This is a good time to display the advantages of tensor notation.

Notation used above
$\mathrm{X}_{\mathrm{u}}$ and $\mathrm{X}_{\mathrm{v}}$
$\mathrm{W}=\mathrm{a} \mathrm{X}_{\mathrm{u}}+\mathrm{b} \mathrm{X}_{\mathrm{v}}$
$\quad \begin{aligned} & \text { Tensor notation } \\ & \mathrm{X}_{, 1} \\ & \mathrm{~W}=\mathrm{w}^{1} \mathrm{X}_{, 2} \\ & \mathrm{~W}=\mathrm{X}_{, 1}+\mathrm{w}^{2} \mathrm{X}_{, 2} \\ & \quad=\mathrm{w}^{\mathrm{i}} \mathrm{X}_{, \mathrm{i}} \\ & \mathrm{Y}=\mathrm{y}^{\mathrm{i}} \mathrm{X}_{, \mathrm{i}} \\ & \mathrm{Y}\left(\mathrm{w}^{1}\right) \text { and } \mathrm{Y}\left(\mathrm{w}^{2}\right) \\ & \left.\mathrm{D}_{\mathrm{Y}} \mathrm{W} \text { (or } \nabla_{\mathrm{Y}} \mathrm{W}\right)\end{aligned}$.

Formula for covariant derivative

$$
\begin{aligned}
\mathrm{DW} / \mathrm{dt}= & \left(\mathrm{a}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{u}} \\
& +\left(\mathrm{b}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{v}} .
\end{aligned}
$$

Same formula in tensor notation

$$
\mathrm{D}_{\mathrm{Y}} \mathrm{~W}=\left(\mathrm{Y}\left(\mathrm{w}^{\mathrm{k}}\right)+\mathrm{w}^{\mathrm{i}} \Gamma_{\mathrm{ij}}^{\mathrm{k}} \mathrm{y}^{\mathrm{j}}\right) \mathrm{X}_{\mathrm{k}} .
$$

## Parallel vector fields and parallel transport.

Let $S$ be a regular surface in $R^{3}$, and $\alpha: I \rightarrow S$ a smooth curve in S . A vector field W along $\alpha$ is a choice of tangent vector $\mathrm{W}(\mathrm{t}) \in \mathrm{T}_{\alpha(\mathrm{t})} \mathrm{S}$ for each $\mathrm{t} \in \mathrm{I}$.

This vector field is smooth if we can write

$$
\mathrm{W}(\mathrm{t})=\mathrm{a}(\mathrm{t}) \mathrm{X}_{\mathrm{u}}+\mathrm{b}(\mathrm{t}) \mathrm{X}_{\mathrm{v}}
$$

in local coordinates, with $a(t)$ and $b(t)$ smooth fns of $t$.
Problem 1. Check that this definition of smoothness of a vector field along $\alpha$ is independent of the choice of local coordinates for $S$.

Example. The velocity vector field $\alpha^{\prime}(\mathrm{t})$ is an example of a smooth vector field along $\alpha$.

If W is a smooth vector field along the smooth curve $\alpha$ on $S$, then the expression
$\mathrm{DW} / \mathrm{dt}=\left(\mathrm{a}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{u}}$

$$
+\left(\mathrm{b}^{\prime}+\mathrm{a} \Gamma_{11}^{2} \mathrm{u}^{\prime}+\mathrm{a} \Gamma_{12}^{2} \mathrm{v}^{\prime}+\mathrm{b} \Gamma_{21}^{2} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{v}}
$$

is well-defined and is called the covariant derivative of W along $\alpha$. As before, $\mathrm{DW} / \mathrm{dt}$ is simply the orthogonal projection of $\mathrm{dW} / \mathrm{dt}$ onto $\mathrm{T}_{\mathrm{p}} \mathrm{S}$.

Example. Let $\alpha$ be a smooth curve on the regular surface $S$, with velocity vector field $\alpha^{\prime}(\mathrm{t})$. The covariant derivative $D \alpha^{\prime} / \mathrm{dt}$ is the portion of the acceleration $\mathrm{d} \alpha^{\prime} / \mathrm{dt}=\alpha^{\prime \prime}(\mathrm{t})$ which is tangent to $S$.

Definition. A smooth vector field W defined along a smooth curve $\alpha: I \rightarrow S$ is said to be parallel if

$$
\mathrm{DW} / \mathrm{dt}=0 \text { for all } \mathrm{t} \epsilon \mathrm{I}
$$

Problem 2. Show that a vector field W defined along a curve $\alpha$ in the plane $R^{2}$ is parallel along $\alpha$ if and only if W is constant.

Problem 3. Let V and W be parallel vector fields along a curve $\alpha: \mathrm{I} \rightarrow \mathrm{S}$. Show that the inner product $\langle\mathrm{V}, \mathrm{W}\rangle$ is constant along $\alpha$. Conclude that the lengths $|\mathrm{V}|$ and $|\mathrm{W}|$ are also constant along $\alpha$.

Problem 4. Let $\alpha: I \rightarrow S^{2}$ parametrize a great circle at constant speed. Show that the velocity field $\alpha^{\prime}$ is parallel along $\alpha$.

Proposition. Let $\alpha: I \rightarrow$ S be a smooth curve on the regular surface S . Let $\mathrm{W}_{0}$ be an arbitrary tangent vector to S at $\alpha\left(\mathrm{t}_{0}\right)$. Then there is a unique parallel vector field $\mathrm{W}(\mathrm{t})$ along $\alpha$ with $\mathrm{W}\left(\mathrm{t}_{0}\right)=\mathrm{W}_{\mathbf{0}}$.

Proof. Working in local coordinates $\mathrm{X}: \mathrm{U} \rightarrow \mathrm{S}$, we can write $\alpha(\mathrm{t})=\mathrm{X}(\mathrm{u}(\mathrm{t}), \mathrm{v}(\mathrm{t}))$. Let

$$
W(t)=a(t) X_{u}+b(t) X_{v}
$$

be the vector field we seek.

Then, since

$$
\begin{aligned}
\mathrm{DW} / \mathrm{dt}= & \left(\mathrm{a}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{u}} \\
& +\left(\mathrm{b}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{v}},
\end{aligned}
$$

the condition that $\mathrm{W}(\mathrm{t})$ be parallel along $\alpha$ is that

$$
\begin{aligned}
& \mathrm{a}^{\prime}+\mathrm{a} \Gamma_{11}^{1} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{1}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{22} \mathrm{v}^{\prime}=0 \\
& \mathrm{~b}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma_{22}^{2} \mathrm{v}^{\prime}=0
\end{aligned}
$$

This is a system of two first order linear ODEs for the unknown functions $a(t)$ and $b(t)$. By standard theorems, a solution exists and is unique, with given initial condition $W_{0}=a\left(t_{0}\right) X_{u}+b\left(t_{0}\right) X_{v}$.

Remark. This proposition allows us to talk about parallel transport of a given tangent vector $\mathrm{W}_{0} \in \mathrm{~T}_{\mathrm{p}} \mathrm{S}$ along a curve $\alpha$ on $S$ which passes through $p$.

Problem 5. Let $\alpha$ be a smooth curve on $S$ connecting the points p and q . Show that parallel transport along $\alpha$ is an isometry from $T_{p} S$ to $T_{q} S$.

Problem 6. Show that if two surfaces are tangent along a common curve $\alpha$, then parallel transport along $\alpha$ is the same for both surfaces.

Problem. Explain how to carry out parallel transport along piecewise smooth curves.

## Geodesics.

Definition. Let $S$ be a regular surface in $\mathrm{R}^{3}$. A smooth curve $\gamma: \mathrm{I} \rightarrow \mathrm{S}$ is called a geodesic if the field of its tangent vectors $\gamma^{\prime}(\mathrm{t})$ is parallel along $\gamma$, that is, if

$$
\mathrm{D} \gamma^{\prime} / \mathrm{dt}=0 .
$$

Note that we can also write this equation as

$$
\mathrm{D}_{\gamma^{\prime}} \gamma^{\prime}=0 \quad \text { or } \quad \nabla_{\gamma^{\prime}} \gamma^{\prime}=0 .
$$

## Remarks.

- The geodesics on the plane $\mathrm{R}^{2}$ are just the straight lines, travelled at constant speed.
- Every geodesic on a surface is travelled at constant speed.
- A straight line which lies on a surface is automatically a geodesic.
- A smooth curve on a surface is a geodesic if and only if its acceleration vector is normal to the surface.
- The geodesics on a round sphere are the great circles.

Problem 7. (a) Find as many geodesics as you can on the right circular cylinder $x^{2}+y^{2}=1$ in $R^{3}$.
(b) Observe that there can be infinitely many geodesics connecting two given points on this cylinder.

Next we want to define the geodesic curvature of a curve on a regular surface. Before doing that, let's recall how we defined curvature of curves in $\mathrm{R}^{3}$ and $\mathrm{R}^{2}$.

If $\alpha: I \rightarrow R^{3}$ is a smooth curve parametrized by arc length, we defined the curvature of $\alpha$ at $s$ to be the real number $\kappa(s)=\left|\alpha^{\prime \prime}(s)\right|$. There is no way to give a sign to the curvature of a curve in $\mathrm{R}^{3}$.

But if $\alpha: I \rightarrow R^{2}$ is a smooth plane curve parametrized by arc length, we can give a sign to its curvature as follows.

Orient the plane $\mathrm{R}^{2}$. Let $\mathrm{T}(\mathrm{s})=\alpha^{\prime}(\mathrm{s})$ be the unit tangent vector to the curve at $\alpha(\mathrm{s})$. Let $\mathrm{N}(\mathrm{s})$ to be the unit vector normal to $\mathrm{T}(\mathrm{s})$ such that the ordered O.N. basis $\mathrm{T}(\mathrm{s}), \mathrm{N}(\mathrm{s})$ agrees with the chosen orientation of $\mathrm{R}^{2}$.

Then define $\kappa(\mathrm{s})=\left\langle\alpha^{\prime \prime}(\mathrm{s}), \mathrm{N}(\mathrm{s})\right\rangle$.
With the usual orientation of $\mathrm{R}^{2}$, positive curvature indicates the curve is bending to the left as you go ahead; negative indicates bending to the right.

We can do the same thing on an oriented regular surface $S$ in $\mathrm{R}^{3}$, as follows.

Let $\alpha: I \rightarrow S$ be a smooth curve on $S$, parametrized by arc length. Let $T(s)=\alpha^{\prime}(s)$ be the unit tangent vector to the curve at $\alpha(\mathrm{s})$. Let $\mathrm{M}(\mathrm{s})$ be the unit vector at $\alpha(\mathrm{s})$ which is tangent to the surface $S$ but orthogonal to $T(s)$, chosen so that the ordered O.N. basis $\mathrm{T}(\mathrm{s}), \mathrm{M}(\mathrm{s})$ agrees with the chosen orientation of $T_{\alpha(s)} S$.

If we have already chosen a unit surface normal N for our surface $S$, then we can simply let $\mathrm{M}(\mathrm{s})=\mathrm{N}(\alpha(\mathrm{s})) \times \mathrm{T}(\mathrm{s})$. That way, the ordered O.N. basis $\mathrm{T}(\mathrm{s}), \mathrm{M}(\mathrm{s}), \mathrm{N}(\alpha(\mathrm{s}))$ agrees with the orientation of $\mathrm{R}^{3}$.

Now we define the geodesic curvature of the curve $\alpha$ at the point $\alpha(\mathrm{s})$ to be

$$
\kappa_{\mathrm{g}}(\mathrm{~s})=\left\langle\alpha^{\prime \prime}(\mathrm{s}), \mathrm{M}(\mathrm{~s})\right\rangle .
$$

Note that
$\left\langle\alpha^{\prime \prime}(\mathrm{s}), \mathrm{M}(\mathrm{s})\right\rangle=\left\langle\mathrm{d} \alpha^{\prime} / \mathrm{d} \mathrm{s}, \mathrm{M}(\mathrm{s})\right\rangle=\left\langle\mathrm{D} \alpha^{\prime} / \mathrm{ds}, \mathrm{M}(\mathrm{s})\right\rangle$.
Thus a smooth curve $\gamma: I \rightarrow S$ parametrized by arc length is a geodesic if and only if its geodesic curvature is zero.

Problem 8. Show that the geodesic curvature of the curve $\alpha: I \rightarrow S$ at the point $\alpha(s)$ is the same as the ordinary curvature at that point of the plane curve obtained by projecting $\alpha$ orthogonally onto the tangent plane $T_{\alpha(s)} S$.

Solution. Shift parameters so that $s=0$ at the point in question, assume that $\alpha(0)$ is at the origin of our coordinate system, and that the unit tangent vector T to this curve, the normal M within the surface, and the surface normal N line up with the $\mathrm{x}, \mathrm{y}$ and z axes.

Assume that s is an arc length parameter along our curve.

Write $\quad \alpha(s)=(x(s), y(s), z(s))$.
Then $\alpha^{\prime}(0)=\left(x^{\prime}(0), y^{\prime}(0), z^{\prime}(0)\right)=(1,0,0)$
and $\quad \alpha "(0)=\left(x "(0), y^{\prime \prime}(0), z^{\prime \prime}(0)\right)=(0, b, c)$.
When $s=0$, the curvature of this curve in 3-space is

$$
\left|\alpha^{\prime \prime}(0)\right|=\left(b^{2}+c^{2}\right)^{1 / 2}
$$

while its geodesic curvature on the surface is $b$.

Projecting our curve onto the tangent plane to the surface at the given point yields the curve

$$
\beta(s)=(x(s), y(s), 0),
$$

where s is no longer an arc length parameter.
The unsigned curvature of $\beta$ at the given point is

$$
\begin{aligned}
\kappa(0) & =\left|\beta^{\prime}(0) \times \beta^{\prime \prime}(0)\right| /\left|\beta^{\prime}(0)\right|^{3} \\
& =|(1,0,0) \times(0, b, 0)| / 1^{3}=|\mathrm{b}| .
\end{aligned}
$$

Its signed curvature there is b , the same as the geodesic curvature of $\alpha$ on our surface, completing the argument.

Problem 9. Let $\alpha: I \rightarrow S$ be a smooth curve on the regular surface $S$ in $R^{3}$.

Let $\kappa(s)$ be the ordinary curvature of the curve $\alpha$ in $R^{3}$, let $\kappa_{\mathrm{g}}(\mathrm{s})$ be its geodesic curvature on the surface S , and let $k_{n}(s)$ be the normal curvature of the surface $S$ at the point $\alpha(\mathrm{s})$ in the direction $\alpha^{\prime}(\mathrm{s})$. Show that

$$
\kappa(\mathrm{s})^{2}=\kappa_{\mathrm{g}}(\mathrm{~s})^{2}+\mathrm{k}_{\mathrm{n}}(\mathrm{~s})^{2} .
$$

Check this when $\alpha$ is a small circle on a round sphere.

Solution. Let's do the example first on a sphere of radius R , as shown below.


The small circle shown has radius $\mathrm{r}=\mathrm{R} \sin \theta$, and when parametrized by arc length is given by

$$
\alpha(\mathrm{s})=(\mathrm{r} \cos (\mathrm{~s} / \mathrm{r}), \mathrm{r} \sin (\mathrm{~s} / \mathrm{r}), \mathrm{R} \cos \theta) .
$$

Then

$$
\begin{aligned}
& \alpha^{\prime}(\mathrm{s})=(-\sin (\mathrm{s} / \mathrm{r}), \cos (\mathrm{s} / \mathrm{r}), 0) \quad \text { and } \\
& \alpha^{\prime \prime}(\mathrm{s})=(-(1 / \mathrm{r}) \cos (\mathrm{s} / \mathrm{r}),-(1 / \mathrm{r}) \sin (\mathrm{s} / \mathrm{r}), 0)
\end{aligned}
$$

The curvature of this small circle is $\kappa=1 / \mathrm{r}=1 /(\mathrm{R} \sin \theta)$.
Its geodesic curvature is

$$
\kappa_{\mathrm{g}}=\alpha^{\prime \prime} \cdot \mathrm{M}=(1 / \mathrm{r}) \cos \theta=\cos \theta /(\mathrm{R} \sin \theta)
$$

The normal curvature of our surface is $\kappa_{n}=1 / R$.

To check that $\kappa^{2}=\kappa_{\mathrm{g}}{ }^{2}+\kappa_{\mathrm{n}}{ }^{2}$, we write

$$
\begin{aligned}
\kappa_{\mathrm{g}}^{2}+\kappa_{\mathrm{n}}^{2} & =\left(\cos ^{2} \theta /\left(\mathrm{R}^{2} \sin ^{2} \theta\right)\right)+1 / \mathrm{R}^{2} \\
& =\left(\cos ^{2} \theta+\sin ^{2} \theta\right) /\left(\mathrm{R}^{2} \sin ^{2} \theta\right) \\
& =1 /\left(\mathrm{R}^{2} \sin ^{2} \theta\right)=\kappa^{2}
\end{aligned}
$$

as desired.

The proof of the formula

$$
\kappa(\mathrm{s})^{2}=\kappa_{\mathrm{g}}(\mathrm{~s})^{2}+\mathrm{k}_{\mathrm{n}}(\mathrm{~s})^{2}
$$

is an application of Meusnier's Theorem.


The curve $\alpha$ shown above passes through the point p on the surface $S$ with tangent vector $T$ and principal normal $N_{\alpha}$. Its curvature there is $\kappa$.

The orthormal frame T, M, N consists of the tangent vector T, an orthogonal vector $M$ still tangent to $S$, and the surface normal N .

The plane spanned by T and N cuts the surface along the curve $\alpha_{o}$, whose curvature $k_{n}$ is the normal curvature of the surface $S$ at $p$.

By Meusnier's Theorem, the curvature $\kappa$ of the curve $\alpha$ at $p$ is related to the normal curvature $\mathrm{k}_{\mathrm{n}}$ by the formula

$$
\kappa=\mathrm{k}_{\mathrm{n}} / \cos \theta,
$$

where $\theta$ is the angle between $\mathrm{N}_{\alpha}$ and N , as shown in the figure above.

Assuming $\alpha$ is parametrized by arc length, its geodesic curvature $\kappa_{g}$ at $p$ is by definition
$\kappa_{g}=\left\langle\alpha^{\prime \prime}, \mathrm{M}\right\rangle=\left\langle\kappa \mathrm{N}_{\alpha}, \mathrm{M}\right\rangle=\kappa \cos (\pi / 2-\theta)=\kappa \sin \theta$.
Then

$$
\kappa_{\mathrm{g}}^{2}+\mathrm{k}_{\mathrm{n}}^{2}=\kappa^{2} \sin ^{2} \theta+\kappa^{2} \cos ^{2} \theta=\kappa^{2},
$$

as desired.

THEOREM (Existence and uniqueness of geodesics). Let S be a regular surface in $\mathrm{R}^{3}$, p a point on S , and $\mathrm{W} \neq 0$ a tangent vector to S at p . Then there is an $\varepsilon>0$ and a unique geodesic $\gamma:(-\varepsilon, \varepsilon) \rightarrow S$ such that

$$
\gamma(0)=p \quad \text { and } \quad \gamma^{\prime}(0)=\mathbf{W} .
$$

Proof. Using local coordinates $X: U \rightarrow S$, let us write the geodesic to be found as $\gamma(\mathrm{t})=\mathrm{X}(\mathrm{u}(\mathrm{t}), \mathrm{v}(\mathrm{t}))$. Then $\gamma^{\prime}(\mathrm{t})=\mathrm{X}_{\mathrm{u}} \mathrm{u}^{\prime}(\mathrm{t})+\mathrm{X}_{\mathrm{v}} \mathrm{v}^{\prime}(\mathrm{t})$.

In our now familiar formula for the covariant derivative,

$$
\begin{aligned}
\mathrm{DW} / \mathrm{dt}= & \left(\mathrm{a}^{\prime}+\mathrm{a} \Gamma_{11}^{1} \mathrm{u}^{\prime}+\mathrm{a} \Gamma_{12}^{1} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{1}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{u}} \\
& +\left(\mathrm{b}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{11} \mathrm{u}^{\prime}+\mathrm{a} \Gamma^{2}{ }_{12} \mathrm{v}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{21} \mathrm{u}^{\prime}+\mathrm{b} \Gamma^{2}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{v}},
\end{aligned}
$$

the role of the vector $W=a X_{u}+b X_{v}$ will be played by $\gamma^{\prime}(t)=X_{u} u^{\prime}(t)+X_{v} v^{\prime}(t)$, and the role of the vector $\mathrm{Y}=\mathrm{u}^{\prime} \mathrm{X}_{\mathrm{u}}+\mathrm{v}^{\prime} \mathrm{X}_{\mathrm{v}}$ will also be played by $\gamma^{\prime}(\mathrm{t})$.

In other words,

$$
\mathrm{a}=\mathrm{u}^{\prime}, \quad \mathrm{b}=\mathrm{v}^{\prime}, \quad \mathrm{a}^{\prime}=\mathrm{u}^{\prime \prime} \quad \text { and } \quad \mathrm{b}^{\prime}=\mathrm{v}^{\prime \prime} .
$$

Thus
$\mathrm{DW} / \mathrm{dt}=\mathrm{D} \gamma^{\prime} / \mathrm{dt}=\mathrm{D}_{\gamma^{\prime}} \gamma^{\prime}$

$$
\begin{aligned}
= & \left(\mathrm{u}^{\prime}+\mathrm{u}^{\prime} \Gamma_{11}^{1} \mathrm{u}^{\prime}+\mathrm{u}^{\prime} \Gamma_{12}^{1} \mathrm{v}^{\prime}+\mathrm{v}^{\prime} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{v}^{\prime} \Gamma_{22}^{1} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{u}} \\
& +\left(\mathrm{v}^{\prime \prime}+\mathrm{u}^{\prime} \Gamma_{11}^{2} \mathrm{u}^{\prime}+\mathrm{u}^{\prime} \Gamma^{2}{ }_{12} \mathrm{v}^{\prime}+\mathrm{v}^{\prime} \Gamma^{2}{ }_{21} \mathrm{u}^{\prime}+\mathrm{v}^{\prime} \Gamma^{2}{ }_{22} \mathrm{v}^{\prime}\right) \mathrm{X}_{\mathrm{v}} .
\end{aligned}
$$

So the system of ODEs to be satisfied by a geodesic is

$$
\begin{aligned}
& \mathrm{u}^{\prime \prime}+\mathrm{u}^{\prime} \Gamma_{11}^{1} \mathrm{u}^{\prime}+\mathrm{u}^{\prime} \Gamma_{12}^{1} \mathrm{v}^{\prime}+\mathrm{v}^{\prime} \Gamma_{21}^{1} \mathrm{u}^{\prime}+\mathrm{v}^{\prime} \Gamma_{22}^{1} \mathrm{v}^{\prime}=0, \\
& \mathrm{v}^{\prime \prime}+\mathrm{u}^{\prime} \Gamma_{11}^{2} \mathrm{u}^{\prime}+\mathrm{u}^{\prime} \Gamma_{12}^{2} \mathrm{v}^{\prime}+\mathrm{v}^{\prime} \Gamma_{21}^{2} \mathrm{u}^{\prime}+\mathrm{v}^{\prime} \Gamma_{22}^{2} \mathrm{v}^{\prime}=0 .
\end{aligned}
$$

The standard existence and uniqueness theorem for such systems of ODEs promises us a unique solution $\mathrm{u}=\mathrm{u}(\mathrm{t}), \mathrm{v}=\mathrm{v}(\mathrm{t})$ defined on some interval $(-\varepsilon, \varepsilon)$ and satisfying the initial conditions

$$
\begin{gathered}
\mathrm{u}(0)=\mathrm{u}_{0}, \mathrm{v}(0)=\mathrm{v}_{0}, \mathrm{u}^{\prime}(0)=\mathrm{u}_{0}^{\prime} \text { and } \mathrm{v}^{\prime}(0)=\mathrm{v}_{0}^{\prime}, \\
\text { where } \mathrm{X}\left(\mathrm{u}_{0}, \mathrm{v}_{0}\right)=\mathrm{p} \text { and } \mathrm{dX}\left(\mathrm{u}_{0}^{\prime}, \mathrm{v}_{0}{ }^{\prime}\right)=\mathrm{W} .
\end{gathered}
$$

This completes the proof.
Problem 10. Show that in tensor notation, the two equations for a geodesic $\gamma(\mathrm{t})=\mathrm{X}\left(\mathrm{u}^{1}(\mathrm{t}), \mathrm{u}^{2}(\mathrm{t})\right)$ are

$$
\left(\mathrm{u}^{\mathrm{k}}\right)^{\prime \prime}+\left(\mathrm{u}^{\mathrm{i}}\right)^{\prime} \Gamma_{\mathrm{ij}}^{\mathrm{k}}\left(\mathrm{u}^{\mathrm{j}}\right)^{\prime}=0, \quad \text { for } \mathrm{k}=1,2 .
$$

## Example - Geodesics on a surface of revolution.

Consider the surface of revolution parametrized by

$$
X(u, v)=(f(v) \cos u, f(v) \sin u, g(v))
$$


$X_{u}=(-f \sin u, f \cos u, 0)$
$X_{v}=\left(f^{\prime} \cos u, f^{\prime} \sin u, g^{\prime}\right)$

$$
\begin{aligned}
& E=\left\langle X_{u}, X_{u}\right\rangle=f^{2} \\
& F=\left\langle X_{u}, X_{v}\right\rangle=0 \\
& \mathrm{G}=\left\langle\mathrm{X}_{\mathrm{v}}, \mathrm{X}_{\mathrm{v}}\right\rangle=\mathrm{f}^{\prime 2}+\mathrm{g}^{\prime 2} \\
& \Gamma_{11}^{1}=0 \\
& \Gamma_{11}^{2}=-\mathrm{ff}^{\prime} /\left(\mathrm{f}^{\prime 2}+\mathrm{g}^{\prime 2}\right) \\
& \Gamma_{12}^{1}=\mathrm{f}^{\prime} / \mathrm{f} \quad \Gamma^{2}{ }_{12}=0 \\
& \Gamma_{22}^{1}=0 \quad \Gamma_{22}^{2}=\left(f^{\prime} f^{\prime \prime}+g^{\prime} g^{\prime \prime}\right) /\left(f^{\prime 2}+g^{\prime 2}\right)
\end{aligned}
$$

The geodesic equations are
(1) $\mathrm{u}^{\prime \prime}+\mathrm{u}^{\prime} \Gamma^{1}{ }_{11} \mathrm{u}^{\prime}+\mathrm{u}^{\prime} \Gamma^{1}{ }_{12} \mathrm{v}^{\prime}+\mathrm{v}^{\prime} \Gamma^{1}{ }_{21} \mathrm{u}^{\prime}+\mathrm{v}^{\prime} \Gamma^{1}{ }_{22} \mathrm{v}^{\prime}=0$
(2) $\mathrm{v}^{\prime \prime}+\mathrm{u}^{\prime} \Gamma^{2}{ }_{11} \mathrm{u}^{\prime}+\mathrm{u}^{\prime} \Gamma^{2}{ }_{12} \mathrm{v}^{\prime}+\mathrm{v}^{\prime} \Gamma^{2}{ }_{21} \mathrm{u}^{\prime}+\mathrm{v}^{\prime} \Gamma^{2}{ }_{22} \mathrm{v}^{\prime}=0$.

Inserting the actual values for the Christoffel symbols gives
(1) $u^{\prime \prime}+2\left(f^{\prime} / f\right) u^{\prime} v^{\prime}=0$
(2) $\mathrm{v}^{\prime \prime}+\left(-\mathrm{f} \mathrm{f}^{\prime} /\left(\mathrm{f}^{\prime 2}+\mathrm{g}^{\prime 2}\right)\right) \mathrm{u}^{\prime 2}$

$$
+\left(\left(\mathrm{f}^{\prime} \mathrm{f}^{\prime \prime}+\mathrm{g}^{\prime} \mathrm{g}^{\prime \prime}\right) /\left(\mathrm{f}^{\prime 2}+\mathrm{g}^{\prime 2}\right)\right) \mathrm{v}^{\prime 2}=0
$$

## Caution about the notation:

$f^{\prime}=d f / d v \quad f^{\prime \prime}=d^{2} f / d v^{2}$ and likewise for $g$, but
$u^{\prime}=d u / d t \quad u^{\prime \prime}=d^{2} u / d t^{2}$ and likewise for $v$.

Problem 11. Check that $u=$ constant and $v=v(t)$ is a solution of the geodesic equations for some choice of $v(t)$.

Hint. Equation (1) above is automatically satisfied, and equation (2) simply determines $v(t)$ so that the curve is travelled at constant speed.

Problem 12. Show that the curve $X(u(t), v(t))$ on our surface of revolution is travelled at constant speed if and only if

$$
\begin{align*}
u^{\prime \prime} f^{2} u^{\prime} & +v^{\prime \prime}\left(f^{\prime 2}+g^{\prime 2}\right) v^{\prime}  \tag{3}\\
& +\left(f^{\prime} f^{\prime \prime}+g^{\prime} g^{\prime \prime}\right) v^{\prime 3}+f^{\prime} u^{\prime 2} v^{\prime}=0
\end{align*}
$$

Problem 13. Show that equations (1) and (2) together imply equation (3).

Problem 14. Show that if $\mathrm{f}^{\prime}\left(\mathrm{v}_{0}\right)=0$, then the circle $\mathrm{u}=\mathrm{ct}$ and $\mathrm{v}=\mathrm{v}_{0}$ satisfies equations (1) and (2), and is hence a geodesic.

Problem 15. Show that if $v^{\prime} \neq 0$, then equations (1) and (3) together imply equation (2). So to get a geodesic, just satisfy equation (1) and make sure you travel at constant speed.

The issue now is to interpret equation
(1) $u^{\prime \prime}+2\left(f\right.$ '/f) $u^{\prime} v^{\prime}=0$.

This equation implies that

$$
\begin{aligned}
\left(\mathrm{f}^{2} \mathrm{u}^{\prime}\right)^{\prime} & =\mathrm{f}^{2} \mathrm{u}^{\prime \prime}+2 \mathrm{ff}^{\prime} \mathrm{v}^{\prime} \mathrm{u}^{\prime} \\
& =\mathrm{f}^{2}\left(\mathrm{u}^{\prime \prime}+2\left(\mathrm{f}^{\prime} / \mathrm{f}\right) \mathrm{u}^{\prime} \mathrm{v}^{\prime}\right)=0
\end{aligned}
$$

which tells us that $\quad f^{2} u^{\prime}=$ constant.

To see the meaning of this, imagine that we travel at constant speed c along the curve $\mathrm{X}(\mathrm{u}(\mathrm{t}), \mathrm{v}(\mathrm{t}))$ on our surface of revolution.

Let $\alpha(\mathrm{t})$ denote the angle that our curve makes with the horizontal circle on the surface through the given point.

Then on the one hand,

$$
\left\langle\mathrm{X}_{\mathrm{u}} \mathrm{u}^{\prime}+\mathrm{X}_{\mathrm{v}} \mathrm{v}^{\prime}, \mathrm{X}_{\mathrm{u}}\right\rangle=\left\langle\mathrm{X}_{\mathrm{u}}, \mathrm{X}_{\mathrm{u}}\right\rangle \mathrm{u}^{\prime}=\mathrm{f}^{2} \mathrm{u}^{\prime}
$$

while on the other hand, this inner product equals

$$
\left|\mathrm{X}_{\mathrm{u}} \mathrm{u}^{\prime}+\mathrm{X}_{\mathrm{v}} \mathrm{v}^{\prime}\right|\left|\mathrm{X}_{\mathrm{u}}\right| \cos \alpha=\mathrm{c} \mathrm{f} \cos \alpha .
$$

So the equation $\mathrm{f}^{2} \mathrm{u}^{\prime}=$ constant is equivalent to

$$
\begin{equation*}
\mathrm{f} \cos \alpha=\text { constant } \tag{4}
\end{equation*}
$$



CLAIRAUT'S THEOREM. Geodesics on the surface of revolution $\mathrm{X}(\mathrm{u}, \mathrm{v})=(\mathrm{f}(\mathrm{v}) \cos \mathrm{u}, \mathrm{f}(\mathrm{v}) \sin \mathrm{u}, \mathrm{g}(\mathrm{v}))$ are characterized by the equation
$\mathrm{f} \cos \alpha=$ constant.



## Comments.

- The value of $\mathrm{f} \cos \alpha$ is constant along a given geodesic, but different geodesics may have different constants.
- If we consider all geodesics through a given point

$$
\mathrm{X}\left(\mathrm{u}_{0}, \mathrm{v}_{0}\right)=\left(\mathrm{f}\left(\mathrm{v}_{0}\right) \cos \mathrm{u}_{0}, \mathrm{f}\left(\mathrm{v}_{0}\right) \sin \mathrm{u}_{0}, \mathrm{~g}\left(\mathrm{v}_{0}\right)\right)
$$

on the surface, then

$$
-\mathrm{f}\left(\mathrm{v}_{0}\right) \leq \text { constant } \leq \mathrm{f}\left(\mathrm{v}_{0}\right)
$$

The extreme constants $-f\left(v_{0}\right)$ and $f\left(\mathrm{v}_{0}\right)$ correspond to geodesics through $\mathrm{X}\left(\mathrm{u}_{0}, \mathrm{v}_{0}\right)$ which at that point are tangent to the horizontal circle.

The constant 0 corresponds to the vertical geodesic through $\mathrm{X}\left(\mathrm{u}_{0}, \mathrm{v}_{0}\right)$, which is simply the profile curve $\mathrm{u}=\mathrm{u}_{0}$.

- Traveling along a given geodesic, as the surface moves farther away from the z -axis, the geodesic becomes more vertical.

