

Context: where Lie algebras  
meet field theory

Model situation:

$U_g$  as "edge currents" for the  
2d gauge theory known  
as topological BF theory

Recall: char 0 field  $k = \mathbb{R}$  or  $\mathbb{C}$

Every assoc. algebra has an  
"underlying" algebra

$\text{Alg}_{\text{Ass}} \xrightarrow{\text{"forget"}} \text{Alg}_{\text{Lie}}$

$(A, \star) \longmapsto (A, [\cdot, \cdot])$

$$[a, b] = a \star b - b \star a$$



There's a "dual" construction:  
the left adjoint

$$\begin{array}{ccc}
 \text{Alg}_{\text{Lie}} & \xrightarrow{U} & \text{Alg}_{\text{Ass}} \\
 \mathfrak{g} & \xrightarrow{\quad} & U\mathfrak{g} = \text{Tens}(\mathfrak{g}) / \langle xoy - yox \rangle \\
 & \nearrow & = \langle [x, y] \rangle \\
 & \text{universal} & \\
 & \text{enveloping} & \\
 & \text{algebra} & 
 \end{array}$$

This is obviously important in algebra  
but it also shows up in  
quantum mechanics:

$\mathfrak{g} \subset$  quantum mechanical system  
symmetries

$$\rightsquigarrow U\mathfrak{g} \xrightarrow{J} A = \text{algebra of operators for system}$$

A particularly nice situation is  
when  $A$  arises as the



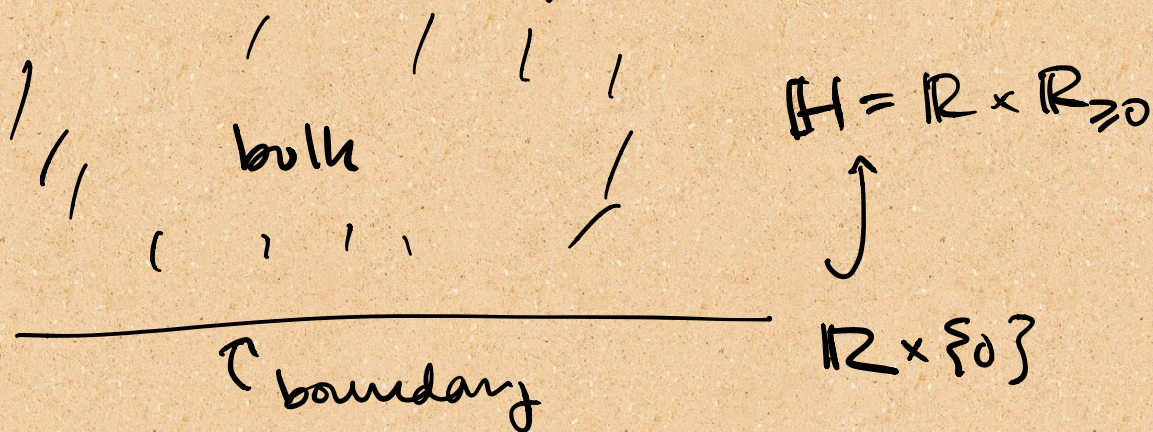




But there's another way that  $U\mathfrak{g}$   
shows up in 2d field theory.

Fix a Lie algebra  $\mathfrak{g}$ .

Consider the "spacetime"



In the bulk, we place topological  
BF theory with fields

$$A \in \Omega^1(\mathbb{H}) \otimes \mathfrak{g}$$

$\uparrow$  connection 1-form

$$B \in \Omega^0(\mathbb{H}) \otimes \mathfrak{g}^*$$

and action functional



$$S_{BF}(A, B) = \int_H (B, F_A)$$

$$F_A = dA + \frac{1}{2} [A, A]$$

The equations of motion are

$$\begin{cases} F_A = 0 \rightsquigarrow \text{flat connections} \\ d_A B = dB + [A, B] \rightsquigarrow B \text{ cov. constant} \end{cases}$$

so

$$\text{Sol} \simeq T^*[\Gamma] \text{Flat}_G$$

cotangent  
direction  
is B

moduli of flat  
 $G$ -connections

For a region  $U$  in the bulk



bulk

the observables  
are



$$\text{Obs}^{\text{cl}}(U) \approx \mathcal{O}(\text{Sol}(U)) \quad \checkmark$$

$$\approx \mathcal{O}(T^*[\text{Flat}_g(U)])$$

We impose the "boundary condition" that  $A|_{\partial H} = 0$  but  $B$  is free.

Hence we're asking that along  $\partial H$ ,  $dB = 0$  so it's constant to a solution  $B \in \mathfrak{g}^*$

$\leadsto$  bulk



$$\text{Obs}^{\text{cl}}(U) \approx \mathcal{O}(\mathfrak{g}^*)$$

$$= \text{Sym}(\mathfrak{g})$$

Claim You can perturbatively quantize this bulk-boundary system so that



or

$$\text{Obs}^{\mathcal{Q}}(V) \underset{\substack{\nearrow \\ \text{bdry} \\ \text{opers}}}{\approx} U_{\mathcal{G}} \quad \text{with } U_{\mathcal{G}} \text{ as "edge currents"}$$

$$\text{Obs}^{\mathcal{Q}}(U) \underset{\substack{\nearrow \\ \text{bulk} \\ \text{opers}}}{\approx} \text{Hoch}^+(U_{\mathcal{G}}, U_{\mathcal{G}})$$

physics

algebra

Remark This claim is part of body of work that grew out of Kontsevich's def. quantization results.



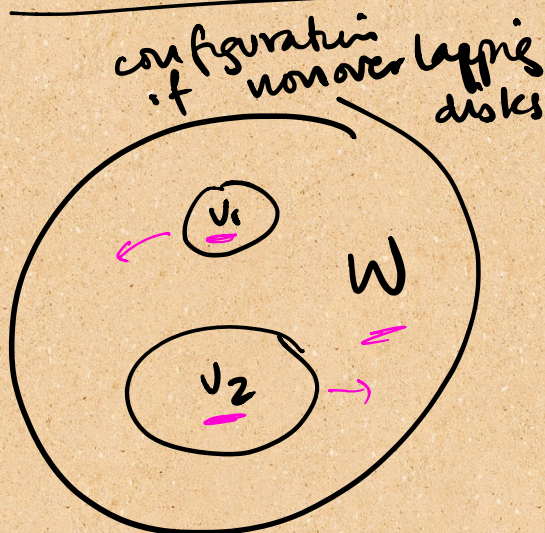
Aside : Eugene Rabinovich  
(graduating student at Berkeley)  
is developing a version of  
Costello's machinery that  
encompasses such bulk-bdry systems  
———— end of model ————

Today: Describe joint work w/  
Ginot, Williams, Zerkalian  
tackling such situations via  
the kind of universal properties  
that Ug has



## Bulk algebras

In the bulk, the observables  $\text{Obs}$  has the structure of an algebra over the little 2-disks operad  $\mathcal{E}_2$



$$\text{Obs}(V_1) \otimes \text{Obs}(V_2)$$



$$\text{Obs}(W)$$

There is a systematic generalization to an operad

$$\mathcal{E}_n = \text{little } n\text{-disks operad}$$

where multiplications are parametrized by configurations of disjoint  $n$ -dim disks in  $\mathbb{R}^n$

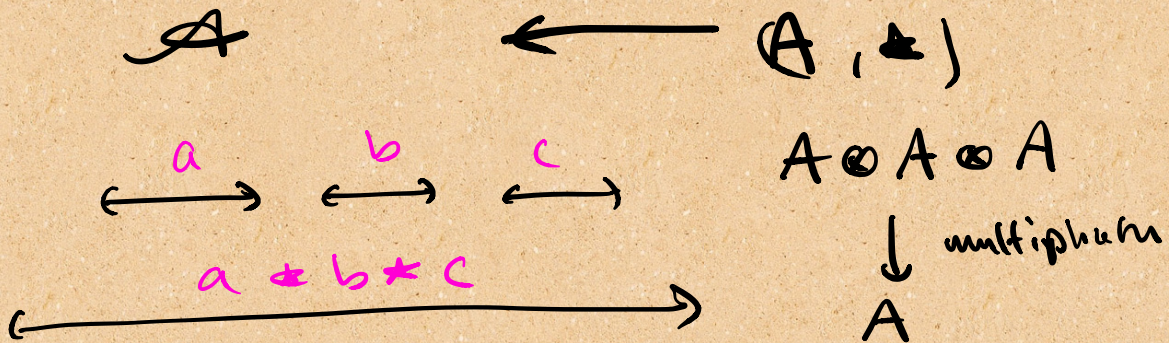


As a special case  $n=1$ , it's helpful to know

$$\underbrace{\text{Alg}_{\mathcal{E}_1}(\text{Ch}_k^{\otimes})}_{\substack{\mathcal{E}_1 \text{ algebras} \\ \text{in chain cplx}}} \cong \underbrace{\text{Alg}_{\text{Ass}}(\text{Ch}_k^{\otimes})}_{\substack{\text{dg associative} \\ \text{algebras} / k}}$$

*∞-categories*

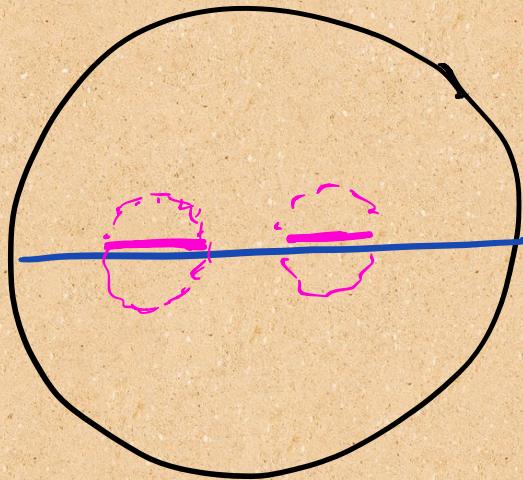
Not too crazy:



You can "forget" from an  $\mathcal{E}_n$ -alg to an  $\mathcal{E}_{n-1}$ -algebra



Ex



$$\begin{array}{c} A \otimes A \\ \downarrow \\ A \end{array}$$

$$\text{Alg}_{\varepsilon_2} \longrightarrow \text{Alg}_{\varepsilon_1}$$

~> hierarchy of "more commutative" algebras

1d

2d

3d

$\text{Alg}_{\varepsilon_1}$

$\leftarrow \text{Alg}_{\varepsilon_2}$

$\leftarrow \text{Alg}_{\varepsilon_3}$

$\leftarrow \dots$

$\leftarrow \text{Alg}_{\varepsilon_\infty}$

12

12

$\text{Alg}_{\text{Ass}}$

$\text{Alg}_{\text{Com}}$

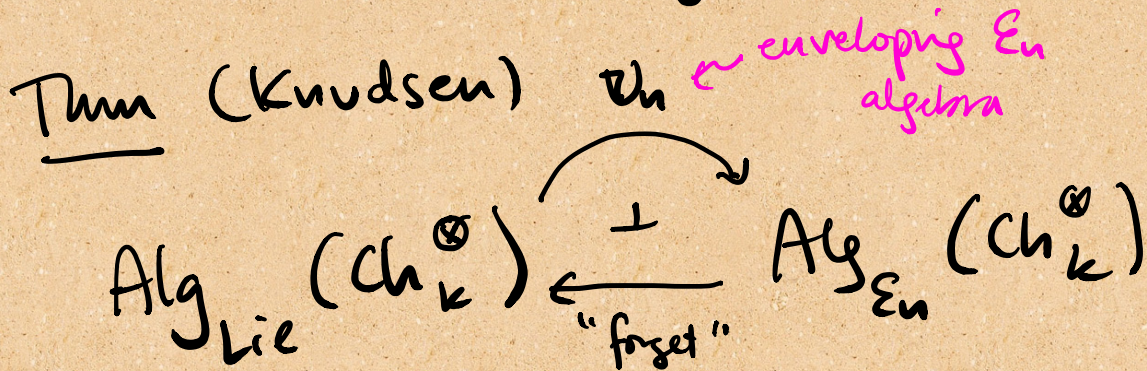
associative

commutative



What about enveloping algebras?

There's a beautiful generalization:



where

- $\mathcal{U}_1 g \cong U_g$  as  $\mathcal{E}_1$ -algs

- there is convenient, explicit model

$$\mathcal{U}_n g(V) \cong C_+^{\text{Lie}}(\underbrace{\Omega_c^*(V) \otimes g}_{\text{de Rham cplx}})$$

$V$  open set  
in  $\mathbb{R}^n$

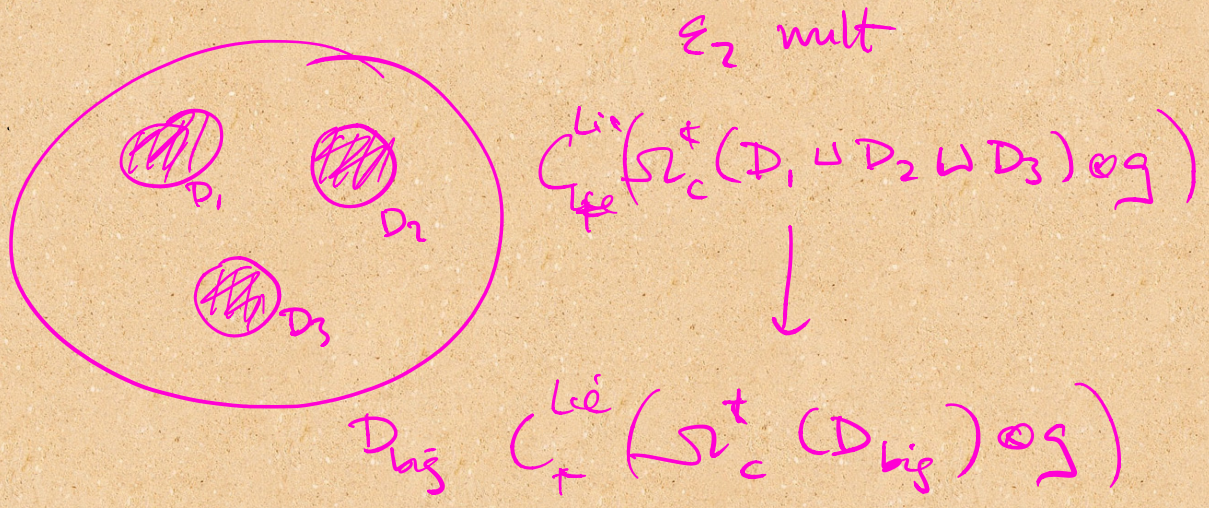
Chesvaly-  
Eilenberg  
chain

compactly  
supp.

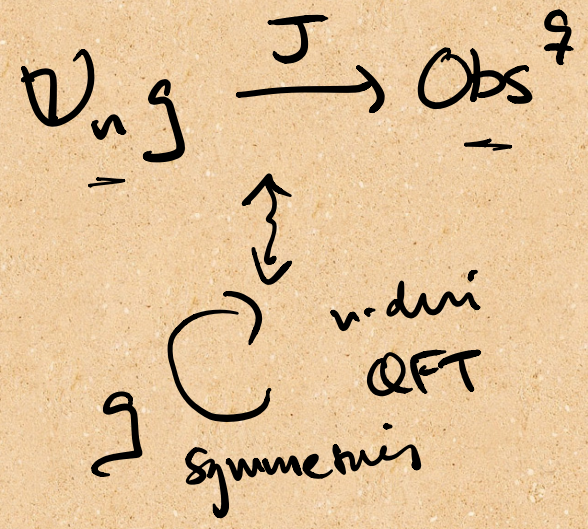
de Rham  
cplx

dg Lie algebra





Runk Kundsen used this to examine  
 config. spaces of wfs but  
 this thing shows up to encode  
 symmetries of n-dim QFTs:

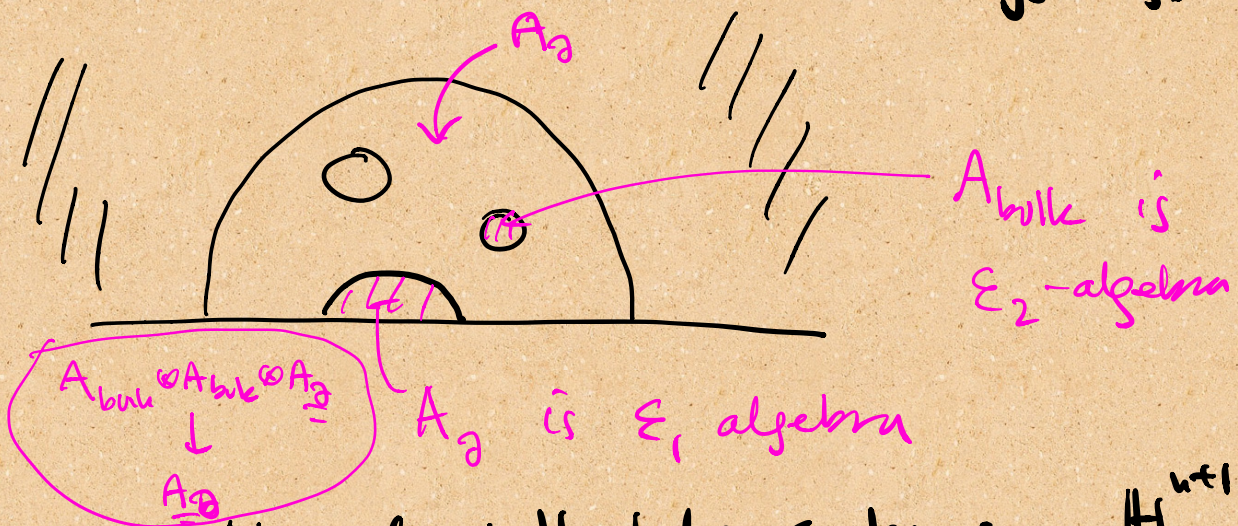


This is explored in Vol 2 of Costello.  
 Tachikawa showed how the



axial anomaly & local index  
fit into this approach.

But let's get back to bulk-bdry  
systems.



Observables of bulk-bdry systems on  $H^{n+1}$   
consist of pairs

$$\begin{array}{ccc} (A_{\text{bulk}} & \circlearrowleft & A_2) \in \text{Alg}_{SC_n} \\ \uparrow & & \uparrow \\ \text{Alg}_{E_{n+1}} & & \text{Alg}_{E_n} \end{array}$$

swiss cheese  
operad

→ configurations of  
disks & half-disks  
in  $H^{n+1}$



Def The  $\mathcal{E}_n$ -center  $\mathcal{Z}_n(A)$  of an  $\mathcal{E}_n$ -algebra  $A$  is the universal  $\mathcal{E}_{n+1}$ -algebra that fits as part of a  $SC_n$ -algebra  $(\mathcal{Z}_n(A), A)$ .

Thm (GGWZ)

For a dg Lie algebra  $\mathfrak{g}$ , the  $\mathcal{E}_n$ -center of  $\mathcal{D}_n \mathfrak{g}$  is modeled by

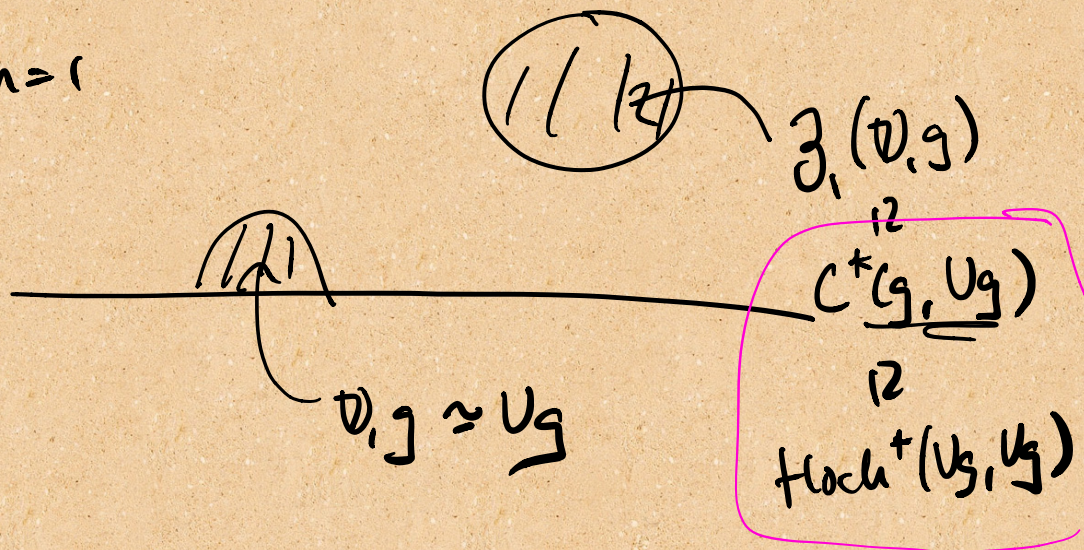
$$\mathcal{Z}_n(\mathcal{D}_n \mathfrak{g})(V) \simeq C_{\text{Lie}}^*(\mathfrak{g}, \mathcal{D}_n \mathfrak{g})$$

$$\simeq C_{\text{Lie}}^*(\underline{\Omega}^* \otimes \mathfrak{g}, C_+^{\text{Lie}}(\underline{\Omega}_c^* \otimes \mathfrak{g}))$$

convenient, explicit model



Ex  $n=1$



### Consequences

① Top. BF theory has a distinguished quantization satisfying  $\text{obs}^{\mathfrak{g}} \simeq \mathcal{Z}_n(V, \mathfrak{g})$

② Lots of fun games looking for generalizations of classic results about  $U\mathfrak{g}$

③  $\int_{M^n} \mathcal{Z}_n(V, \mathfrak{g}) \simeq C_{\text{lie}}^*(\Omega^+(M) \otimes \mathfrak{g}, C_{\tau}^{\text{lie}}(\Omega_c^+(M, \mathfrak{g})))$