

## Solving the Black-Scholes equation

Now we can divide through by  $dt$  to get the *Black-Scholes equation*:

$$\frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + rS \frac{\partial V}{\partial S} - rV = 0.$$

In this equation, we're looking for  $V(S, t)$  and the interest rate  $r$  and the stock's volatility  $\sigma$  are "known" constants. It's interesting that the stock's growth rate  $\mu$  doesn't appear in the equation at all.

Now we (that is, you) need to solve the equation with various "final" conditions at time  $T$ . In particular, we need to do this for  $C$  and  $P$  with the conditions given above.

To derive the solution, the main part of the work is to convert the Black-Scholes equation into the usual heat equation. To do this, you'll have to make three kinds of changes of variable:

- To get the time running in the right direction, you can define a new variable  $\tau = T - t$ . Then  $t = T$  will correspond to  $\tau = 0$ .
- Since it was  $dS/S = d(\log S)$  that satisfied the standard Wiener process that leads to the usual heat equation, it makes sense to define a new variable  $x = \log S$  (natural logarithm). This should get rid of the appearances of the independent variable  $S$  or  $x$  multiplying the various derivatives.
- As we did at the beginning of the course, a substitution of the form  $u = e^{\alpha x + \beta \tau} V$  can be used to get rid of unwanted constants and first-order in  $x$  terms.

With these hints, show that the value of a European call option with strike price  $E$  and expiry time  $T$  is given by:

$$C(S, t) = SF(A_+) - Ee^{-r(T-t)}F(A_-)$$

where  $F(x)$  is the cumulative distribution function for the standard normal distribution:

$$F(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-p^2/2} dp,$$

and the constants  $A_{\pm}$  are given by

$$A_{\pm} = \frac{\log(S/E) + (r \pm \frac{1}{2}\sigma^2)(T-t)}{\sigma\sqrt{T-t}}$$

Here we go – first, if  $\tau = T - t$  then  $\frac{\partial V}{\partial \tau} = -\frac{\partial V}{\partial t}$ , so the equation becomes

$$\frac{\partial V}{\partial \tau} = \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + rS \frac{\partial V}{\partial S} - rV$$

which looks more like the (forward) heat equation.

Next, let  $x = \log S$  (or  $S = e^x$ ). Then

$$\frac{\partial V}{\partial S} = \frac{\partial V}{\partial x} \frac{dx}{dS} = \frac{1}{S} \frac{\partial V}{\partial x}$$

and

$$\frac{\partial^2 V}{\partial S^2} = \frac{1}{S} \frac{\partial}{\partial x} \left( \frac{1}{S} \frac{\partial V}{\partial x} \right) = \frac{1}{S^2} \frac{\partial^2 V}{\partial x^2} - \frac{1}{S^2} \frac{\partial V}{\partial x}$$

(because  $1/S = e^{-x}$  so  $d(1/S)/dx = -e^{-x} = -1/S$ ) and the equation becomes

$$\begin{aligned} \frac{\partial V}{\partial \tau} &= \frac{1}{2} \sigma^2 S^2 \left( \frac{1}{S^2} \frac{\partial^2 V}{\partial x^2} - \frac{1}{S^2} \frac{\partial V}{\partial x} \right) + rS \left( \frac{1}{S} \frac{\partial V}{\partial x} \right) - rV \\ &= \frac{1}{2} \sigma^2 \frac{\partial^2 V}{\partial x^2} + \left( r - \frac{1}{2} \sigma^2 \right) \frac{\partial V}{\partial x} - rV \end{aligned}$$

This last equation is of the form

$$\frac{\partial V}{\partial \tau} = A \frac{\partial^2 V}{\partial x^2} + B \frac{\partial V}{\partial x} + CV$$

for certain constants  $A$ ,  $B$  and  $C$  (with  $A > 0$ ). To avoid dealing with too much algebra, we'll start by proving the following helpful result:

**Lemma.** Suppose  $V(x, \tau)$  satisfies

$$\frac{\partial V}{\partial \tau} = A \frac{\partial^2 V}{\partial x^2} + B \frac{\partial V}{\partial x} + CV$$

with initial condition  $V(x, 0) = f(x)$ , where  $A$ ,  $B$  and  $C$  are constants (with  $A > 0$ ). Then

$$V(x, \tau) = \frac{e^{C\tau}}{\sqrt{4\pi A\tau}} \int_{-\infty}^{\infty} e^{-\frac{1}{2} \left( \frac{y-x-B\tau}{\sqrt{2A\tau}} \right)^2} f(y) dy.$$

*Proof:* Let  $u = e^{\alpha x + \beta \tau} V$ . Then  $V = e^{-(\alpha x + \beta \tau)} u$ , therefore

$$\frac{\partial V}{\partial \tau} = e^{-(\alpha x + \beta \tau)} \left( \frac{\partial u}{\partial \tau} - \beta u \right)$$

and

$$\frac{\partial V}{\partial x} = e^{-(\alpha x + \beta \tau)} \left( \frac{\partial u}{\partial x} - \alpha u \right)$$

and

$$\frac{\partial^2 V}{\partial x^2} = e^{-(\alpha x + \beta \tau)} \left( \frac{\partial^2 u}{\partial x^2} - 2\alpha \frac{\partial u}{\partial x} + \alpha^2 u \right)$$

Putting this into the equation and canceling the  $e^{-(\alpha x + \beta \tau)}$  from both sides gives

$$\frac{\partial u}{\partial \tau} - \beta u = A \left( \frac{\partial^2 u}{\partial x^2} - 2\alpha \frac{\partial u}{\partial x} + \alpha^2 u \right) + B \left( \frac{\partial u}{\partial x} - \alpha u \right) + Cu$$

which we can rearrange as

$$\frac{\partial u}{\partial \tau} = A \frac{\partial^2 u}{\partial x^2} + (B - 2\alpha A) \frac{\partial u}{\partial x} + (C + \beta - \alpha B + \alpha^2 A) u$$

To get rid of the  $\partial u/\partial x$  term, we should choose

$$\alpha = \frac{B}{2A},$$

in which case the equation becomes

$$\frac{\partial u}{\partial \tau} = A \frac{\partial^2 u}{\partial x^2} + \left( C + \beta - \frac{B^2}{4A} \right) u$$

And to get rid of the  $u$  term, we should choose

$$\beta = \frac{B^2}{4A} - C$$

and we arrive at the heat equation

$$\frac{\partial u}{\partial \tau} = A \frac{\partial^2 u}{\partial x^2}$$

with  $k = A$ . And the initial data for  $u$  are

$$u(x, 0) = e^{\alpha x} V(x, 0) = e^{\frac{B}{2A}x} f(x).$$

And the solution of this initial-value problem is

$$u(x, \tau) = \frac{1}{\sqrt{4\pi A\tau}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4A\tau}} e^{\frac{B}{2A}y} f(y) dy$$

Now we need to complete the square in the exponentials, so we calculate

$$\begin{aligned} \frac{(x-y)^2}{4A\tau} - \frac{B}{2A}y &= \frac{1}{4A\tau} (y^2 - (2x + 2B\tau)y + x^2) \\ &= \frac{1}{4A\tau} (y^2 - 2(x + B\tau)y + (x + B\tau)^2 - 2xB\tau - B^2\tau^2) \\ &= \frac{(y - x - B\tau)^2}{4A\tau} - \frac{Bx}{2A} - \frac{B^2\tau}{4A} \\ &= \left( \frac{y - x - B\tau}{2\sqrt{A\tau}} \right)^2 - \alpha x - \beta\tau - C\tau \end{aligned}$$

Therefore

$$u(x, \tau) = \frac{e^{\alpha x + \beta\tau} e^{C\tau}}{\sqrt{4\pi A\tau}} \int_{-\infty}^{\infty} e^{-\left(\frac{y - x - B\tau}{2\sqrt{A\tau}}\right)^2} f(y) dy$$

Multiply by  $e^{-(\alpha x + \beta\tau)}$  and fiddle with the 2's in the exponential and obtain

$$V(x, \tau) = \frac{e^{C\tau}}{\sqrt{4\pi A\tau}} \int_{-\infty}^{\infty} e^{-\frac{1}{2}\left(\frac{y - x - B\tau}{\sqrt{2A\tau}}\right)^2} f(y) dy.$$

This completes the proof of the lemma.

Now we want to apply the lemma with  $A = \frac{1}{2}\sigma^2$ , with  $B = r - \frac{1}{2}\sigma^2$  and with  $C = -r$ . The result is

$$V(x, \tau) = \frac{e^{-r\tau}}{\sigma\sqrt{2\pi\tau}} \int_{-\infty}^{\infty} e^{-\frac{1}{2}\left(\frac{y - x - (r - \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}\right)^2} f(y) dy.$$

Now we apply the initial data. For a call option, we have

$$V(S, t = T) = \begin{cases} S - E & \text{if } S > E \\ 0 & \text{if } S \leq E \end{cases}$$

So this is the data for  $\tau = 0$  and since  $S = e^x$  it makes

$$f(x) = \begin{cases} e^x - E & \text{if } e^x > E \\ 0 & \text{if } e^x \leq E \end{cases}$$

or in other words the “break point” of  $f(x)$  comes at  $\log E$ . So we can rewrite the solution so far as

$$V(x, \tau) = \frac{e^{-r\tau}}{\sigma\sqrt{2\pi\tau}} \int_{\log E}^{\infty} e^{-\frac{1}{2}\left(\frac{y-x-(r-\frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}\right)^2} (e^y - E) dy.$$

Now we make the change of variables:

$$z = \frac{y-x-(r-\frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}} \quad \text{so} \quad dz = \frac{dy}{\sigma\sqrt{\tau}} \quad \text{and} \quad y = x + \left(r - \frac{1}{2}\sigma^2\right)\tau + \sigma\sqrt{\tau}z$$

and rewrite the solution as

$$V(x, \tau) = \frac{e^{-r\tau}}{\sqrt{2\pi}} \int_{\frac{\log E - x - (r - \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}}^{\infty} e^{-\frac{1}{2}z^2} \left( e^{x+(r-\frac{1}{2}\sigma^2)\tau + \sigma\sqrt{\tau}z} - E \right) dz.$$

Almost there! There are two terms. We'll write each one in terms of the cumulative distribution function of the Gaussian, in other words, in terms of the function

$$F(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-p^2/2} dp.$$

It is important to note that

$$F(\infty) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-p^2/2} dp = 1$$

and  $e^{-p^2/2}$  is even, so by the substitution  $p \rightarrow -p$  we can conclude

$$F(-x) = \frac{1}{\sqrt{2\pi}} \int_x^{\infty} e^{-p^2/2} dp,$$

So let's take the second term of  $V$  first, namely

$$\begin{aligned} -E \frac{e^{-r\tau}}{\sqrt{2\pi}} \int_{\frac{\log E - x - (r - \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}}^{\infty} e^{-\frac{1}{2}z^2} dx &= -E e^{-r\tau} F\left(\frac{x - \log E + (r - \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}\right) \\ &= -E e^{-r(T-t)} F\left(\frac{\log(S/E) + (r - \frac{1}{2}\sigma^2)(T-t)}{\sigma\sqrt{T-t}}\right) \\ &= -E e^{-r(T-t)} F(A_-) \end{aligned}$$

using the facts that  $\tau = T - t$ , that  $x = \log S$  and the definition of  $A_-$  given in the problem.

Now we work on the first term:

$$\frac{e^{-r\tau}}{\sqrt{2\pi}} \int_{\frac{\log E - x - (r - \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}}^{\infty} e^{-\frac{1}{2}z^2} e^{x+(r-\frac{1}{2}\sigma^2)\tau + \sigma\sqrt{\tau}z} dz$$

We can factor  $e^x$  out of the integrand as  $S$ . Then we need to complete the square in the rest of the exponent:

$$\begin{aligned} -\frac{1}{2}\left(z^2 - 2\sigma\sqrt{\tau}z - (2r - \sigma^2)\tau\right) &= -\frac{1}{2}\left((z^2 - \sigma\sqrt{\tau})^2 - \sigma^2\tau - (2r - \sigma^2)\tau\right) \\ &= -\frac{1}{2}\left((z - \sigma\sqrt{\tau})^2 - 2r\tau\right) \end{aligned}$$

Using this, we can rewrite the first term as

$$\frac{e^{-r\tau}}{\sqrt{2\pi}} \int_{\frac{\log E - x - (r - \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}}^{\infty} e^{-\frac{1}{2}z^2} e^{x + (r - \frac{1}{2}\sigma^2)\tau + \sigma\sqrt{\tau}z} dz = \frac{S}{\sqrt{2\pi}} \int_{\frac{\log E - x - (r - \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}}^{\infty} e^{-\frac{1}{2}(z - \sigma\sqrt{\tau})^2} dz$$

Now let  $w = z - \sigma\sqrt{\tau}$  so  $dw = dz$  and  $z = w + \sigma\sqrt{\tau}$  and this becomes

$$\frac{S}{\sqrt{2\pi}} \int_{\frac{\log E - x - (r + \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}}^{\infty} e^{-\frac{1}{2}w^2} dw$$

(pay particular attention that  $r - \frac{1}{2}\sigma^2$  changed to  $r + \frac{1}{2}\sigma^2$  in the lower limit of integration because we added  $\sigma\sqrt{\tau}$  in changing from  $z$  to  $w$ ). And finally, using that  $\tau = T - t$ , that  $x = \log S$  and the definition of  $A_+$  given in the problem, we can rewrite this as

$$\begin{aligned} \frac{S}{\sqrt{2\pi}} \int_{\frac{\log E - x - (r + \frac{1}{2}\sigma^2)\tau}{\sigma\sqrt{\tau}}}^{\infty} e^{-\frac{1}{2}w^2} dw &= SF\left(\frac{\log(S/E) + (r + \frac{1}{2}\sigma^2)(T - t)}{\sigma\sqrt{T - t}}\right) \\ &= SF(A_+) \end{aligned}$$

Woohoo!