

LEC 7.

Dimension of c

$$\frac{\Delta u}{(\Delta t)^2} = c^2 \frac{\Delta u}{(\Delta x)^2} \Rightarrow c^2 = \left(\frac{\Delta x}{\Delta t} \right)^2$$

↓
velocity.

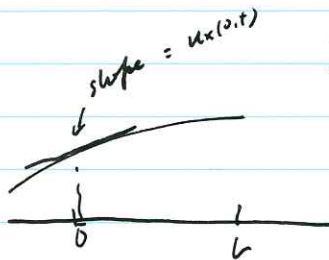
∴ c has the dimension of velocity.

In fact, $f(x-ct)$, $g(x+ct)$ are two travelling wave solutions to $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$.

Boundary condition

1. Dirichlet BCs Fix string at both ends.
 $u(0,t) = y_1(t)$ and $u(L,t) = y_2(t)$.

2. Neumann BCs prescribe the slope of string at end points.



$$\frac{\partial u}{\partial x}(0,t) = y_1(t) \quad \text{and}$$

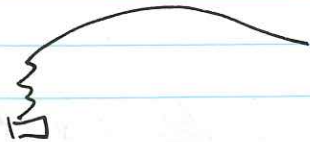
3. Elastic BC

String is attached to mass-spring system with equilibrium position $u_E(t)$, and spring constant k .

$$\text{To } \frac{\partial u}{\partial x}(0,t) = k(u(0,t) - u_E(t)) :$$

nonhomogeneous elastic BC is $u_E(t) \neq 0$.

$u_E = 0 \Rightarrow$ homogeneous elastic BC.



derivation (see book, optional).

Remark: Elastic BC is an analogue of ^{mixed} BC for heat equations (Newton's law of cooling).

Vibrating String with Fixed ends

$$\left. \begin{array}{l} \frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} \quad \text{on } 0 < x < L \\ u(0, t) = u(L, t) = 0 \quad \text{BC} \\ u(x, 0) = f(x) \\ u_t(x, 0) = g(x) \end{array} \right\} \text{IC.}$$

Separation of variables $u(x, t) = \phi(x) h(t)$.

$$\phi(x) \frac{d^2 h}{dt^2} = c^2 \frac{d^2 \phi}{dx^2} h(t)$$

$$\Rightarrow \frac{h''}{c^2 h} = \frac{\phi''}{\phi} = -\lambda$$

eigenvalue problem

$$\left\{ \begin{array}{l} h'' + \lambda c^2 h = 0 \\ \phi'' + \lambda \phi = 0 \\ \phi(0) = \phi(L) = 0 \end{array} \right.$$

time equation $h(t)$.

$$h'' + \lambda c^2 h = 0.$$

$$\lambda > 0. \quad h(t) = C_1 \cos \sqrt{\lambda} c t + C_2 \sin \sqrt{\lambda} c t$$

$$\lambda = 0 \quad h(t) = C_1 + C_2 t$$

$$\lambda < 0 \quad h(t) = C_1 e^{C\sqrt{-\lambda}t} + C_2 e^{-C\sqrt{-\lambda}t}$$

Physically, we expect oscillatory solution in time $\Rightarrow \lambda > 0$.

space equation

$$\left. \begin{array}{l} \phi'' + \lambda \phi = 0 \\ \phi(0) = \phi(L) = 0 \end{array} \right\} \Rightarrow \text{the only non trivial case is when } \lambda > 0.$$

$$\phi(0) = \phi(L) = 0 \quad \Rightarrow \quad \lambda_n = \left(\frac{n\pi}{L}\right)^2, \quad n=1, 2, \dots$$

$$\therefore u(x,t) = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi x}{L} \cos \frac{n\pi ct}{L} + \sum_{n=1}^{\infty} B_n \sin \frac{n\pi x}{L} \sin \frac{n\pi ct}{L}$$

$$\text{IC: } t=0, \quad u(x,0) = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi x}{L} \quad \rightarrow \text{Fourier sine series for } f(x).$$

$$u_t(x,0) = \sum_{n=1}^{\infty} B_n \frac{n\pi c}{L} \sin \frac{n\pi x}{L} = g(x) \quad \rightarrow \text{Fourier sine series for } g(x).$$

Use the orthogonality.

$$A_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$

$$B_n \frac{n\pi c}{L} = \frac{2}{L} \int_0^L g(x) \sin \frac{n\pi x}{L} dx$$

$$\Rightarrow B_n = \frac{2}{n\pi c} \int_0^L g(x) \sin \frac{n\pi x}{L} dx.$$

Note: Heat equation: oscillatory in x , exponential decay in t .

Wave equation: oscillatory in x and t .

$n=1 \Rightarrow$ fundamental frequency.

Other frequencies are multiples of the fundamental

solution $u(x,t)$ is an infinite superposition of terms

normal mode of vibration: $\sin \frac{n\pi x}{L} \left(A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L} \right)$.

Trig alg.

$$A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L} = \sqrt{A_n^2 + B_n^2} \left(\underbrace{\frac{A_n}{\sqrt{A_n^2 + B_n^2}}}_{\sin \theta} \cos \frac{n\pi ct}{L} + \underbrace{\frac{B_n}{\sqrt{A_n^2 + B_n^2}}}_{\cos \theta} \sin \frac{n\pi ct}{L} \right)$$

(use $\sin(\alpha + \beta) = \sin \alpha \cos \beta + \cos \alpha \sin \beta$)

$$= \sqrt{A_n^2 + B_n^2} \sin \left(\frac{n\pi ct}{L} + \theta \right)$$

$\gamma = \sqrt{A_n^2 + B_n^2}$: amplitude. θ : phase shift.
 $\tan \theta = \frac{\sin \theta}{\cos \theta} = \frac{A_n}{B_n}$.

The amplitude $\gamma = \sqrt{A_n^2 + B_n^2}$ determines intensity of the sound.

$\omega = \frac{n\pi c}{L}$: # of oscillations per 2π units of time.
 called ~~also~~ circular frequency.

The sound is an infinite superposition of normal modes,

$$\underbrace{\sqrt{A_n^2 + B_n^2} \sin \left(\frac{n\pi ct}{L} + \theta \right)}_{A(t)} \sin \frac{n\pi x}{L} : \text{nth harmonic.}$$

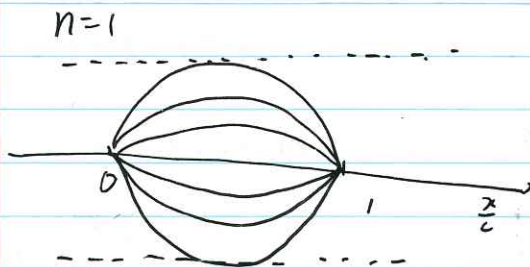
$A(t)$: time-dependent amplitude.

Spatial frequency is $\frac{n\pi}{L}$, while temporal frequency is $\frac{n\pi c}{L}$ which is also called natural frequency.

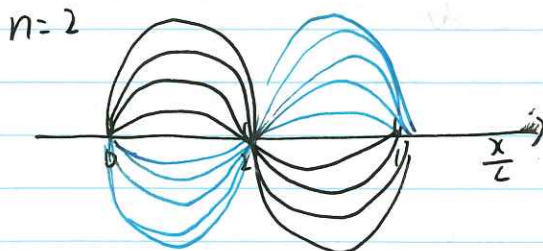
frequencies.

$$w = \frac{\pi c}{L} \text{ depends on } c = \sqrt{\frac{T_0}{\mu_0}}, L.$$

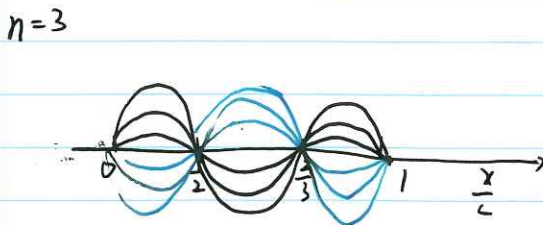
Frequency $w = \frac{\pi c}{L}$ can be increased by increasing c or decreasing L . Higher frequency gives higher pitch. Since μ_0 is fixed, c can be changed by increasing tension T_0 . Length can be clamping down the string.



At time goes, the $\sin\left(\frac{\pi x c t}{L} + \theta\right)$ will vary between -1 and 1 . Since there is no displacement at end points, we have standing waves.



Displacement is zero at $x = \frac{L}{2}$. This point is called a node.



3rd harmonic has 2 nodes at $x = \frac{L}{3}$ and $x = \frac{2L}{3}$.

In general, n th harmonic has $(n-1)$ nodes.

Claim: Any standing waves can be considering the sum of two travelling waves.

trig. identity:

$$\sin \frac{n\pi x}{L} \cos \frac{n\pi ct}{L} = \frac{1}{2} \sin \frac{n\pi}{L} (x-ct) + \frac{1}{2} \sin \frac{n\pi}{L} (x+ct)$$

$$\sin \frac{n\pi x}{L} \sin \frac{n\pi ct}{L} = \frac{1}{2} \cos \frac{n\pi}{L} (x-ct) - \frac{1}{2} \cos \frac{n\pi}{L} (x+ct).$$

Hence $u(x,t)$ is a superposition of standing wave
 \Rightarrow

$$u(x,t) = \underbrace{R(x-ct)}_{\text{wave traveling to the right with speed } c} + \underbrace{L(x+ct)}_{\text{wave traveling to the left with speed } c}.$$

Note: this is true for BC not fixed ends as well.

Recall: $u(x,t) = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi x}{L} \cos \frac{n\pi ct}{L} + \sum_{n=1}^{\infty} B_n \sin \frac{n\pi x}{L} \sin \frac{n\pi ct}{L}$

$$u(x,0) = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi x}{L} = f(x).$$

$$u_t(x,0) = \sum_{n=1}^{\infty} B_n \cdot \frac{n\pi c}{L} \sin \frac{n\pi x}{L} = g(x).$$

Ex: Let $u(x,0) = f(x)$ $u_t(x,0) = 0$.

$$g(x) = 0 \Rightarrow B_n = 0 \quad n=1, 2, 3, \dots$$

$$\therefore u(x,t) = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi x}{L} \cos \frac{n\pi ct}{L}$$

$$= \sum_{n=1}^{\infty} A_n \cdot \frac{1}{2} \left[\sin \frac{n\pi}{L} (x-ct) + \sin \frac{n\pi}{L} (x+ct) \right]$$

$$= \frac{1}{2} \sum_{n=1}^{\infty} A_n \sin \frac{n\pi}{L} (x-ct) + \frac{1}{2} \sum_{n=1}^{\infty} A_n \sin \frac{n\pi}{L} (x+ct).$$

$$= \frac{1}{2} f(x-ct) + \frac{1}{2} f(x+ct).$$

$$\boxed{u(x,t) = \frac{1}{2} f(x-ct) + \frac{1}{2} f(x+ct)}$$